Control of Friction

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Abstract

Spatiotemporal fluctuations in small discrete nonlinear arrays affect the dynamics of the center of mass. We will present numerical evidence indicating that phase synchronization is related to the frictional properties of such sliding atomic scale objects.

We will discuss mechanisms of how the resulting atomic scale friction can be tuned with noise and disorder. We derive a set of two coupled equations describing respectively the motion of the center of mass and the spatial average fluctuations in the coherent mode adopted by the array.

Our analysis of this reduced set of equations indicates that depending on array stiffness and size, quantized jumps in the minimum friction (maximum velocity) of the array occur due to resonant parametric forcing of the particle fluctuations by the center of mass motion. We propose an analytical expression to determine occurrences of these jumps.

Robustness of Friction Mechanisms

Friction is ruled by robust dynamics

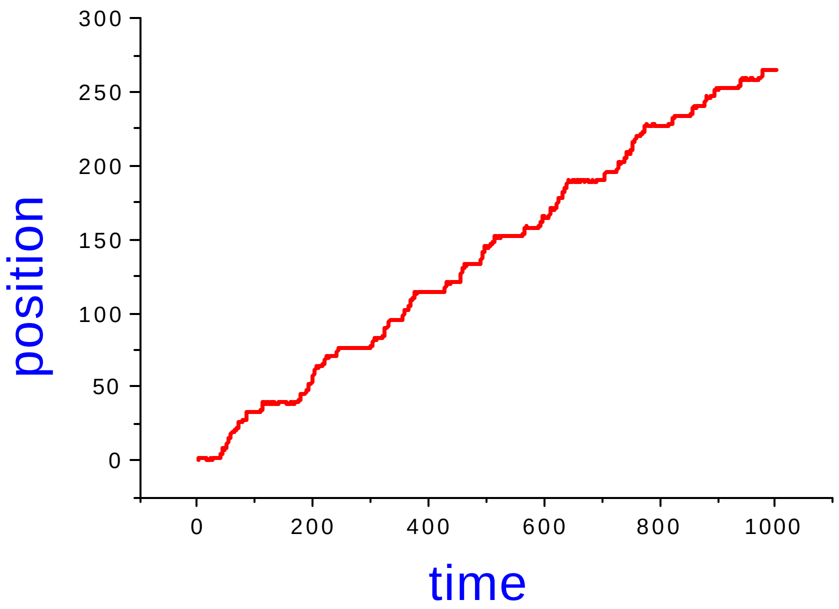
good qualitative agreement between variety of models and types of interaction potentials used for a model

- choice of parameters may be even more important than the choice of a model !!!
- Initial conditions !

Stick-Slip Dynamics

- Has been observed from the nano - to macro scales - from the atomic scale to earthquakes.

Both periodic and chaotic stick-slip dynamics have been observed



Different Regimes of Motion

- Single - particle dynamics

- Collective dynamics

Very limited correlation between particles in array

Propagation of well defined moving structures

High temperature
(high noise)
Large external forcing
Small coupling

Small-medium forcing
Large-intermediate coupling
Reasonable noise/disorder

Collective Dynamics

Understanding collective dynamics is the key issue

It has not been studied before in regard to friction

We have suggested a link from collective motion to friction

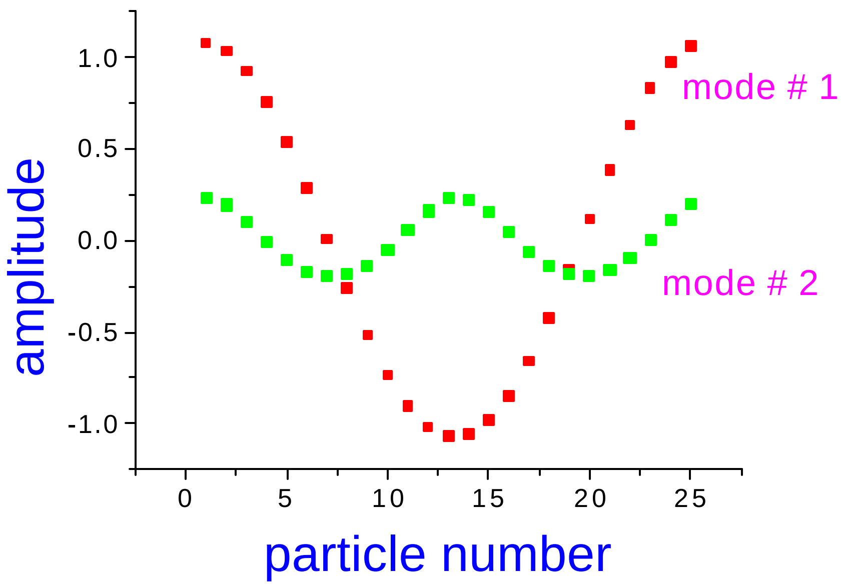
Some of the predictions based on this approach have already been successfully tested experimentally

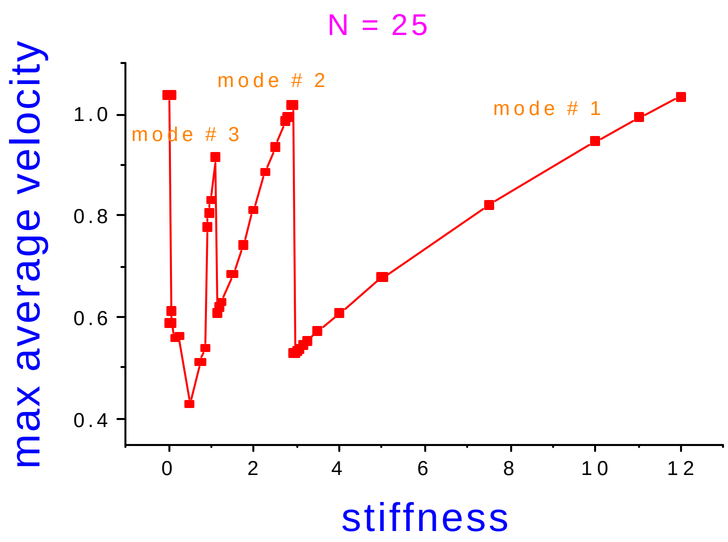
Locking of the Temporal and Spatial Dynamics (Modes)

Small size and confinement

The outcome $\Rightarrow$ Propagation modes

Each mode is characterized by different frictional behavior





Nonlinear Friction Selection

- Simulations using F-K model show that for intermediate to high values of the coupling and small applied force

a series of quantized transitions in the maximum propagation velocity occur.

- It is possible to scale the position at which these maximum velocity jumps occur using the size ***N*** of the array and the coupling κ.
- At low enough values of the coupling a transition back to synchronous motion occurs independent of system size ***N.***

$$\kappa_m(N) \sim (N/m)^2$$

Theoretical Modeling

- Phenomenological models
- F-K-Tomlinson model

$$\mu \ddot{X}_j + \gamma \dot{X}_j = -\partial U / \partial X_j - \partial V / \partial X_j + f_j + \eta_j$$

μ is the mass of the sliding particle

γ is the dissipation coefficient

U is the interaction potential

V is surface potential

f is the external driving force

η is the thermal noise (temperature effect)

Dynamics of Propagating Arrays

We separate the center of mass motion of array from spatiotemporalfluctuations (which only dissipate energy)

$$X_n(t) = X(t) + \delta X_n(t)$$

where $< \delta X_n(t) > = 0$ by construction

Keeping fluctuations small, the center of mass obeys

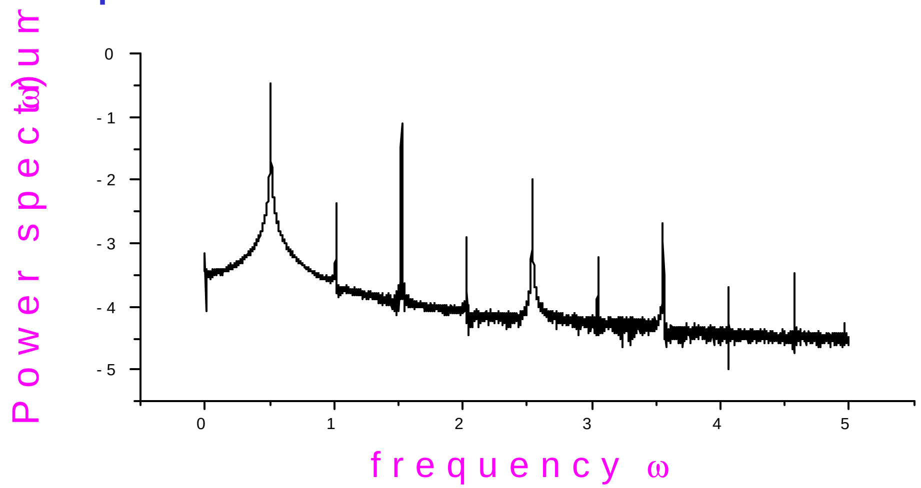
$$\boxed{\ddot{X} + \gamma \dot{X} + \sin(X)[1 - < \delta X_n^2 > /2] = f}$$

The spatiotemporal fluctuations obey

$$\boxed{\delta \ddot{X}_n + \gamma \, \delta \dot{X}_n + C \cos(X) \, \delta X_n = \kappa(\delta X_{n+1} - 2\,\delta X_n + \delta X_{n-1})}$$

Assumptions

We assume that the main mechanism for the energy transfer from the center of mass motion to the spatiotemporal fluctuations in the array is due to a subharmonic parametric resonance.



We have made a self-consistent approximation by replacing nonlinear terms by a quasilinear term.

$$C = \sqrt{1/[1 + 2 < \overline{\delta X_n^2} >]}$$

Resonant Parametric Forcing

We make the Fourier decomposition

$$\delta X_n(t) = \sum_m \delta X_m(t) e^{2\pi imn/N}$$

and equations of motion for the modes

$$\delta \ddot{X}_m + \gamma\, \delta \dot{X}_m + [\Omega_m^2 + C\cos(X)]\delta X_m = 0$$

where $$\Omega_m = 2\sqrt{\kappa}\sin(\pi m/N)$$

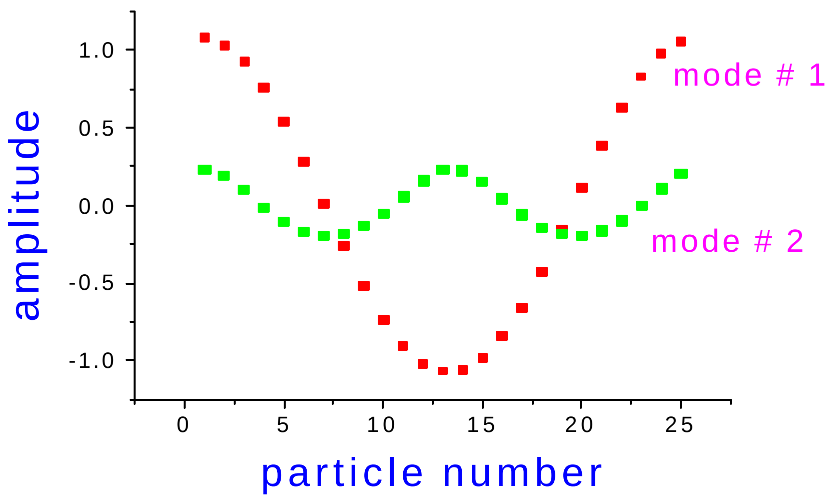
Shows parametric forcing when $\Omega_m = \omega/2$

Spatial Coherence and Mode Selection

If we look for a solution for the ***m***'th mode of the form

$$\delta X_m = b_m \sin(\omega t / 2 + \beta_m)$$

we then find:



Only one mode can exist at a time.

There are N such solutions. Each is spatially coherent with a different center of mass velocity and different amplitude fluctuations.

As the spatial fluctuations b_m increase, phase synchronization decreases, and so the average center of mass velocity decreases.

Velocity of the Center of Mass

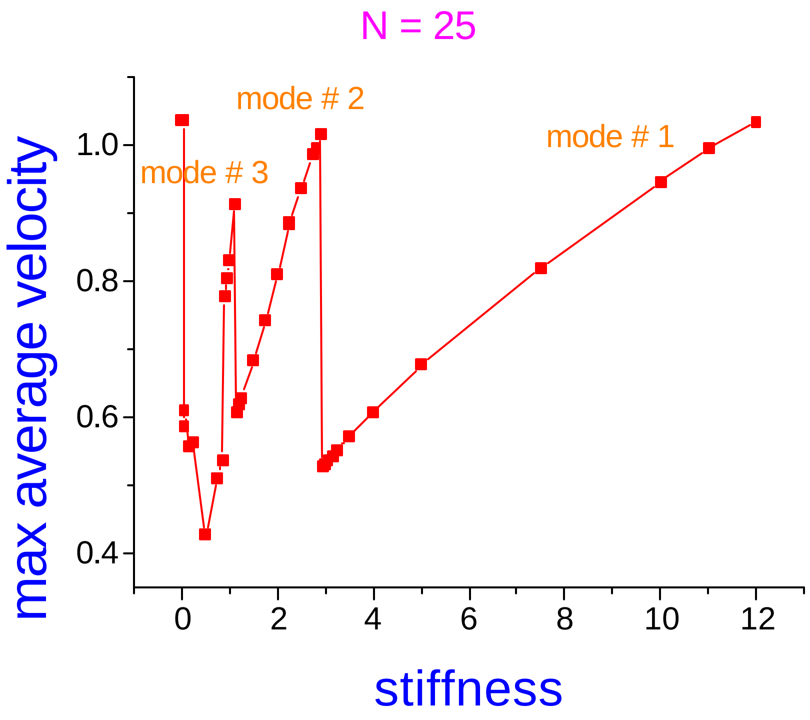
If we look for a solution for the ***m***'th mode of the form

$$\delta X_m = b_m \sin(\omega t/2 + \beta_m)$$

and the center of mass motion is described by

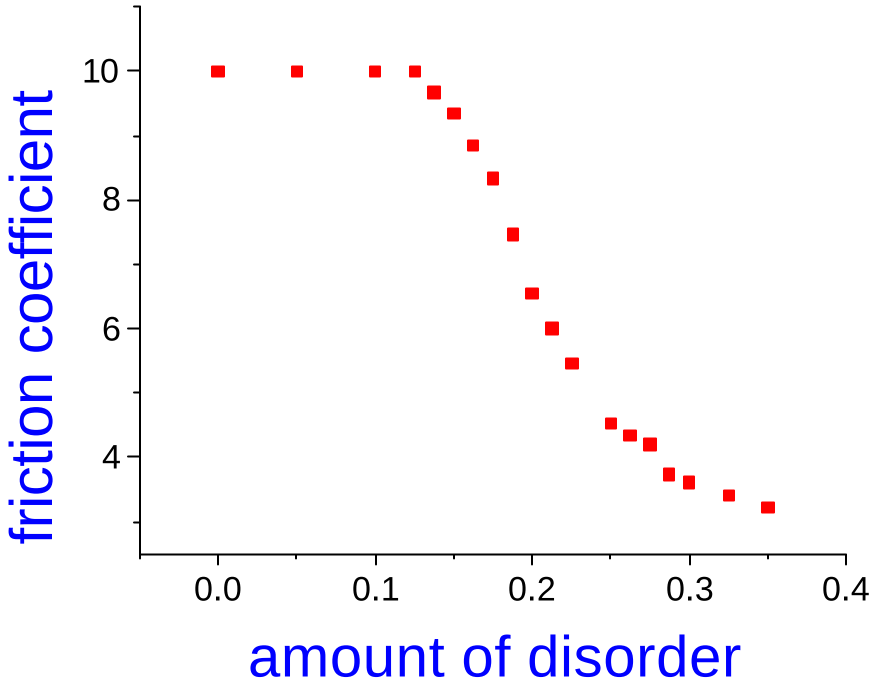
$$X = X_0 + \omega t + B\sin(\omega t)$$

then the velocity of the center of mass is



Sliding on Disordered Substrate

Friction coefficient can be significantly reduced
(by orders of magnitude) when sliding on irregular surfaces



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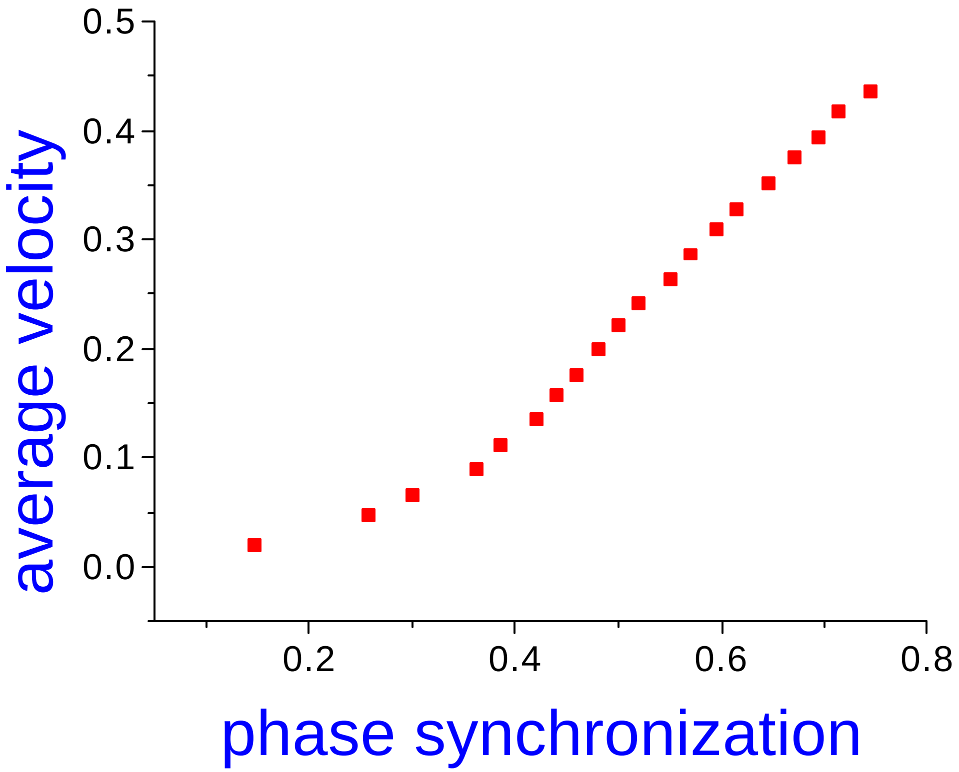
Key Issue ⇒ Phase Synchronization

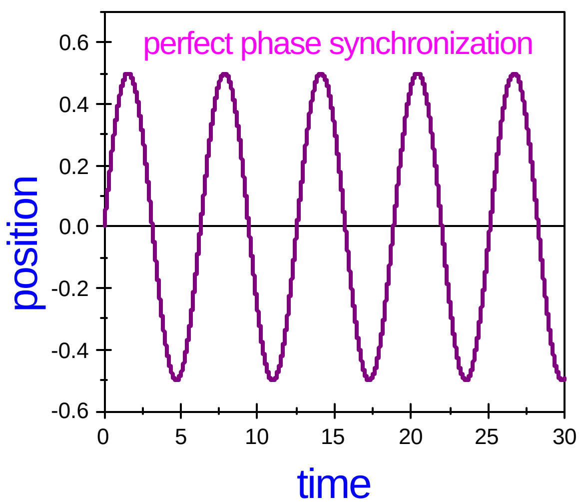


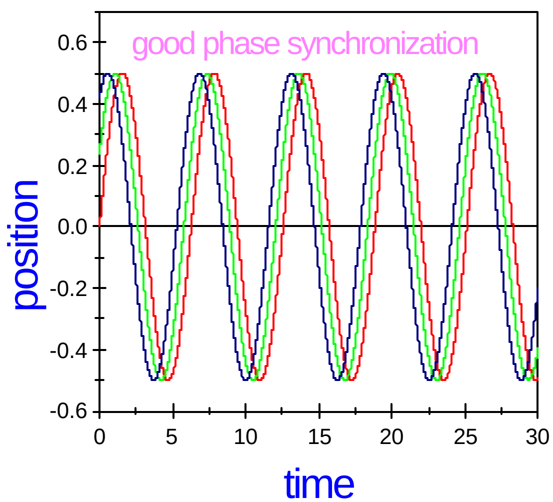
The better the array is phase synchronized - the faster it moves !

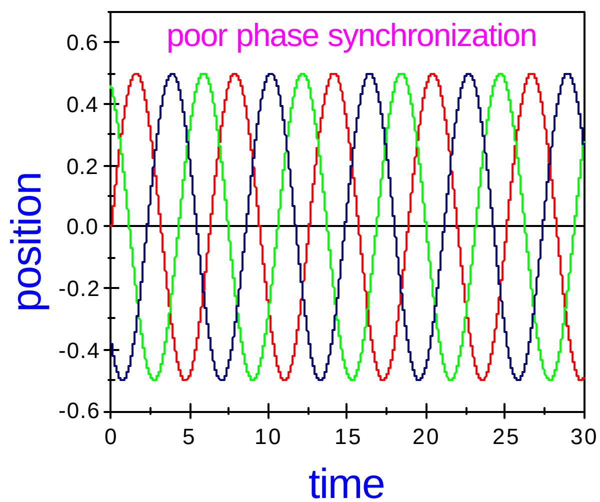
Phase Synchronization

We define phase synchronization as the inverse of the fluctuations σ from the center of mass motion

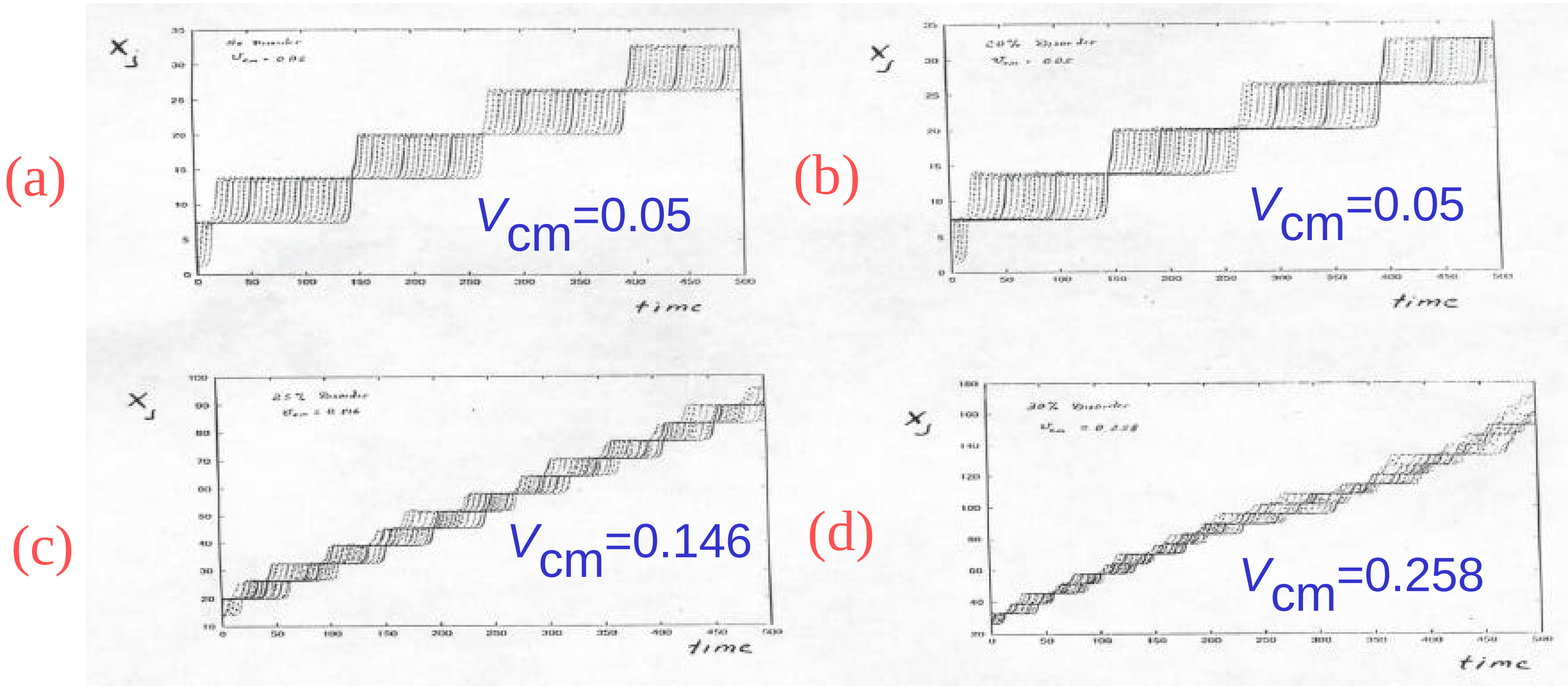
$$\sigma = \lim_{T \to \infty} \frac{1}{T} \int_0^T \frac{\sqrt{\sum_j^N (X_j - X_{av})^2}}{X_{av}}$$



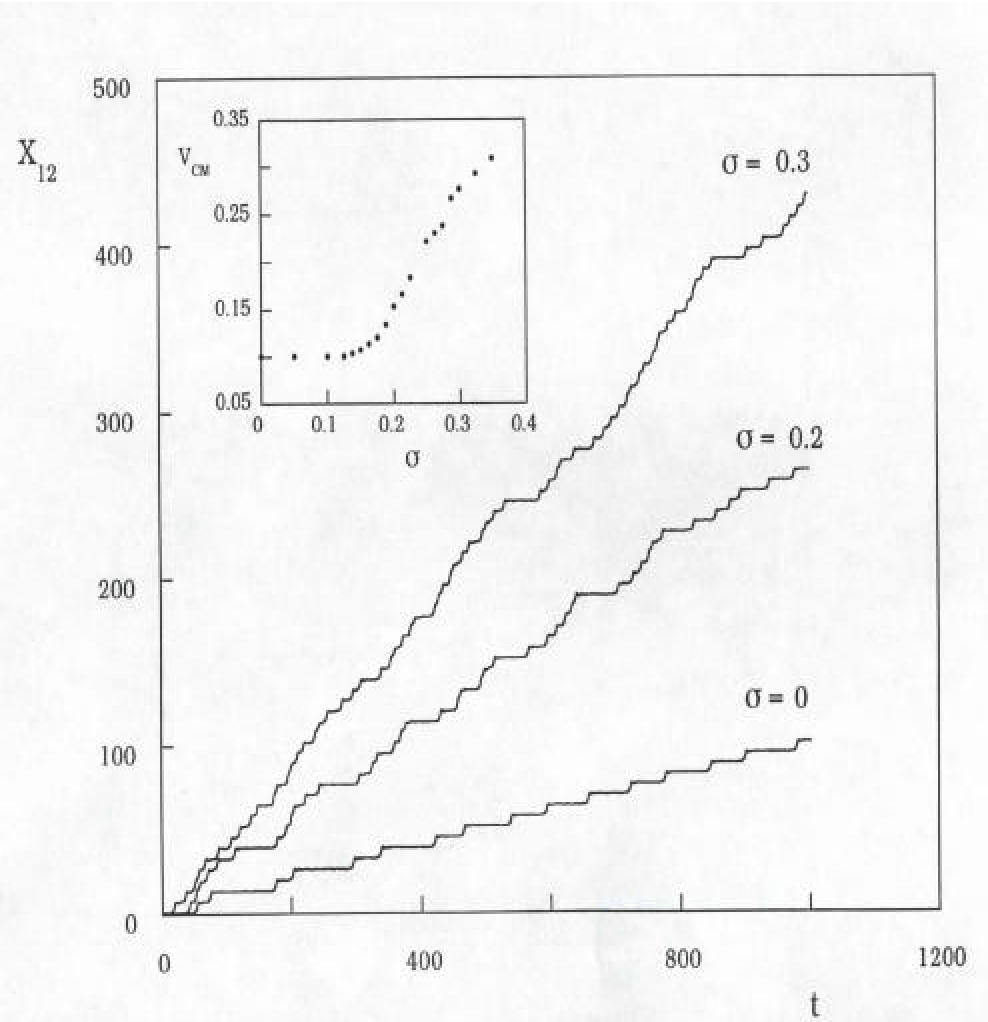




Disorder - Enhanced Synchronization



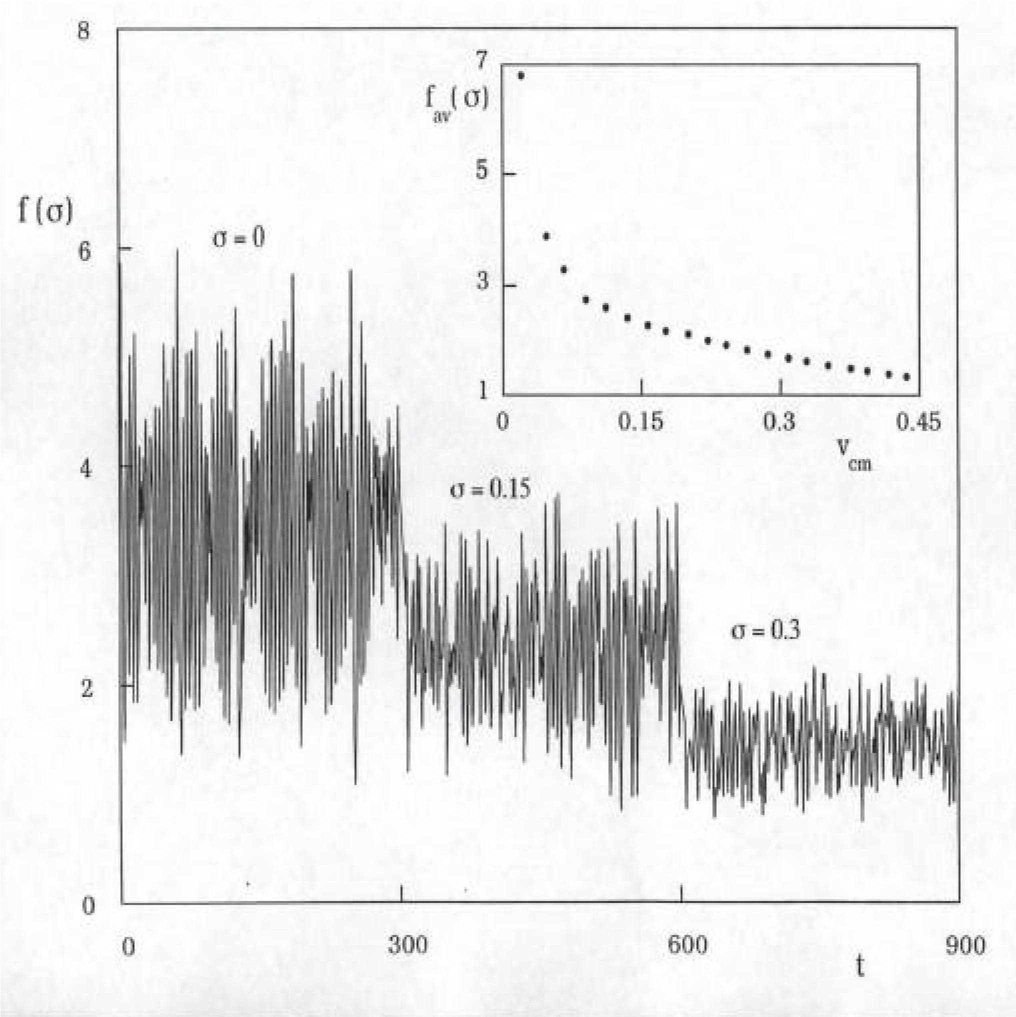
Time series of positions of all the particles in N=25 particle array for:
(a) the identical array; (b) 20% of disorder;
(c) 25 % of disorder; (d) 30 % of disorder



The position of a particle #12
in array as a function of time.

The bottom curve corresponds to the identical array.
The middle curve corresponds to to the arrays with 20% of disorder,
The top curve corresponds to the array with 30% of disorder.

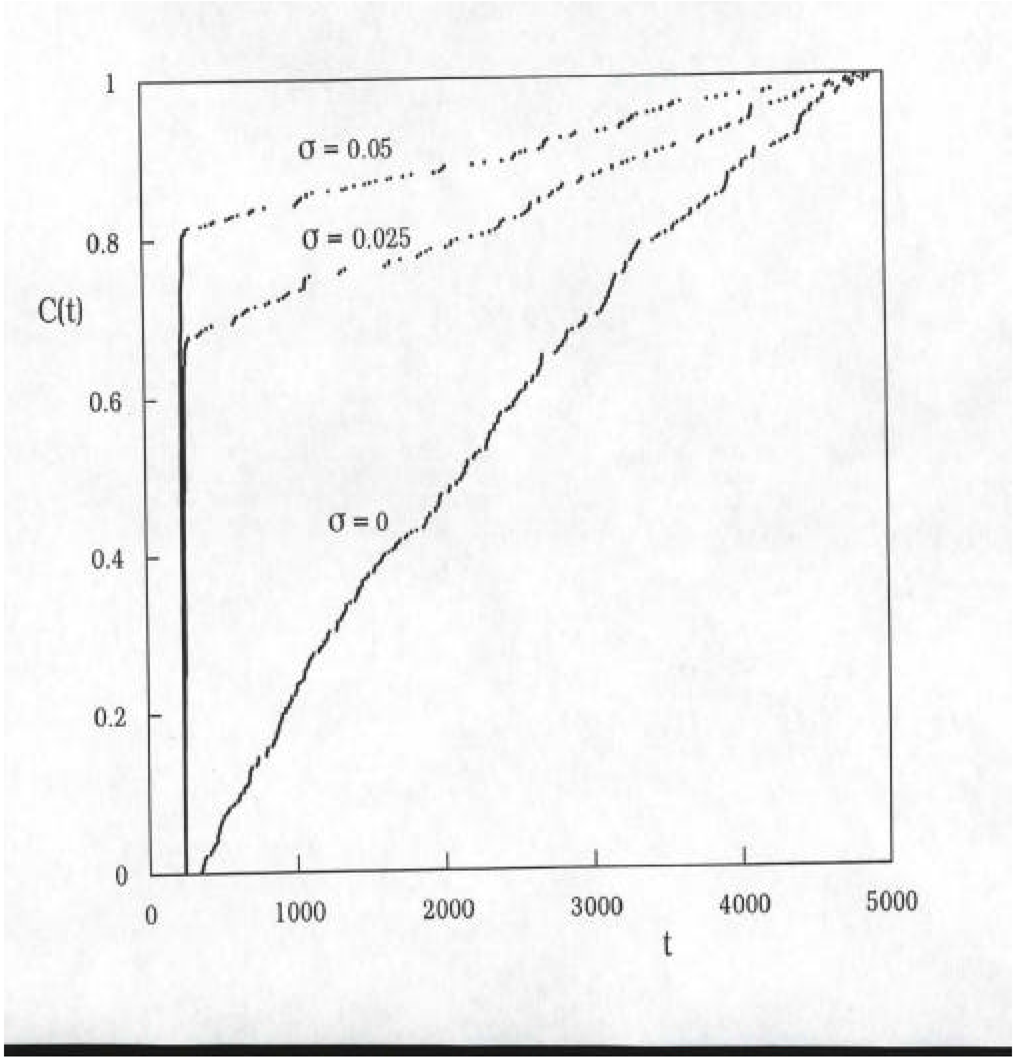
The inset shows the average velocity of the center of mass as a function of the amount of disorder



Time series of the fluctuations from the center of mass f(σ) for different amounts of disorder.

The left-hand part of the plot corresponds to the identical array.
The middle part corresponds to σ=15%.
The right-hand part corresponds to σ=30%.

The inset shows the average fluctuations from the center of mass as the function of the velocity of the center of mass.



Cumulative slip time distribution for the array.

The bottom curve corresponds to the identical array.

The middle curve corresponds to $\sigma = 2.5\%$.

The top curve corresponds to $\sigma = 5\%$.

Summary

Nanoscale arrays can exhibit a variety of modes of motion with different degrees of spatial coherence which affects frictional properties of the array

Energy is transferred between the center of mass motion and the spatiotemporal fluctuations using parametric forcing resulting in mode selection

As a result quantized jumps in the observed friction associated with different spatial mode is possible

Spatial disorder and thermal noise can contribute to increase of the phase synchronization of the array and therefore decrease friction

Spatial disorder and thermal noise can contribute to the depinning process and eliminate stick-slip thus decreasing friction

Publications

- Friction

Phys. Rev. Lett. **83**, 104 (1999)
Phys. Rev. E **59**, R4737 (1999)
Phys. Rev. E **58**, R4057 (1998)
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- Dynamics

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