

**MODELING CARBON REDEPOSITION IN THE TORE SUPRA
ACTIVELY COOLED LIMITER CIEL**

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1. Introduction

The creation of an extrapolable model for carbon erosion and re-deposition processes is needed to evaluate possible tritium retention and diagnostic coating rates in ITER. Tore Supra experiments relevant to this topic have been made, with a high level of heat extraction (>0.6 GJ), for long pulses (t > 4 mins), using the CIEL (Composants Internes et Limiters) actively cooled heat and particle exhaust system [1]. To contribute to the analysis of these experiments, a detailed simulation model for these processes in CIEL has been created, combining the BBQ 3-D Monte Carlo scrape-off layer impurity transport code with core radial impurity transport codes (SANCO / MIST). A BBQ-only model for erosion processes from the CIEL superstructure was compared with spectroscopic measurements from the Limiter Demarrage de CIEL [2, 3]. Here we describe the combined CIEL model (superstructure, leading edge and neutralizer region) required to simulate the re-deposition and total build-up of carbon from a series of long pulse discharges.

2. Model

The model starts from a detailed description of the CIEL geometry using a modification of the CASTEM-2000 thermal analysis code [4]. Using available data for background plasma parameters in the edge/SOL, the 3D incident D⁺ deposition profile is determined. Using, in addition, available information about surface temperature distributions on the CIEL structure and a model for sheath acceleration, the spatial and velocity-space distribution of emitted C fluxes due to the physical, chemical, and radiation enhanced sublimation (RES) processes which are driven by the D⁺ source is calculated. Using these D-related emission distributions, the BBQ Monte Carlo calculation follows the resulting C

evolution in 3-D ($\tau_{\text{equil}} \sim 10$ ms) until final disposition (striking plasma-facing components, or entrance to the core plasma), thus calculating re-deposition and the rate of penetration of carbon into the core plasma region. The resulting charge-distributed C influx to the core plasma then provides the boundary condition for the core radial impurity transport code. The radial code follows the core impurity evolution using this source until equilibration ($\tau_{\text{equil}} \sim 0.1$ s). Subsequently, the calculated C efflux from the radial code constitutes an additional self-sputtered impurity source in BBQ.

While the C influx to the core is drawn from the various CIEL regions, C efflux from the core falls mainly on the super-structure surface. C^0 from the superstructure region is quickly ionized, so that the core influx lies mostly in low ionization stages, while fluxes from the leading edge and neutralizer region dwell for a longer time in the SOL and can attain a higher charge before penetrating the plasma core (defined as $\rho=0.71$ m). After residence in the core for τ_p^+ , however, the core efflux stream has a broad distribution of charges. Each of these species has a different sheath-accelerated energy, and hence self-sputtering yield.

Thus, additional CASTEM maps are created for each impurity ion. Using this information, and the effluxes from the radial transport code, an iteration to account for self-sputtering is made, with the self-sputtered impurities followed by BBQ and the products of self-sputtering re-introduced into the core calculation. The iterative solution allows description of the possible development of slow secular growth in Z_{eff} , which would signal the absence of a true steady state solution. Figure 1 conveys the general outline of the overall scheme with a schematic diagram of the process. Figure 2 shows the convergence of a sample BBQ/MIST iteration.

3. Examples

Calculations using this model are underway for the conditions of a series of Tore Supra shots (29930 – 30093) in which pulse lengths of several of the discharges exceeded 1 minute. Starting with boronization prior to the first discharge, at the end of the sequence a deposited layer of 140 μm was observed at the entrance of the neutralizer [see 5]. In place of

direct measurements of SOL properties, for these sample calculations we assume electron density at the last closed magnetic surface $n_e(\text{LCMS}) = 7 \cdot 10^{18} \text{ m}^{-3}$, $T_e(\text{LCMS}) = 50 \text{ eV}$, and the scrape-off layer decay lengths for density (λ_N) and temperature (λ_T) are 4 cm and 5 cm, respectively. $Z_{\text{eff}} \sim 2$ for long discharges in this shot sequence. Figure 3 shows typical BBQ particle orbits for sources from the super-structure, leading edge and neutralizer regions, along with the (θ, ϕ) distribution of emergent D^+ -sputtered fluxes from these areas. The tendency toward higher charge states for core penetrators originating from the leading edge/neutralizer is demonstrated in Fig. 4, which shows the charge state distribution for C ions entering the core from the various regions. Figure 5 shows the resulting radial distributions of C ion density for the leading edge and neutralizer region sources. High charge state distributions are found deep in the SOL for sources from the leading edge/neutralizer. The use of high power LH in long pulse discharges has resulted in the observation of zones of local high heat deposition on the leading edge and neutralizer regions [1]. The examples discussed here suggest that such a deep SOL source can produce a higher density of more highly ionized C in the SOL than that from the strongly attenuated convective flux to the superstructure, and thus could generate C through self-sputtering at rates consistent with the strong observed build-up. This is a qualitative conjecture, which could be improved with the availability of SOL plasma parameter data from discharges similar to those considered.

Acknowledgement: ORNL. Supported by U.S.DOE Contract DE-AC05-00OR22725.

References

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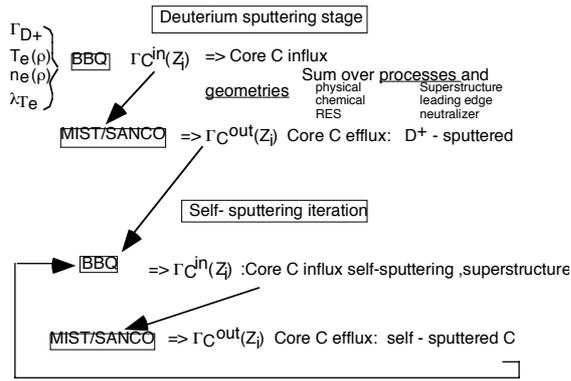


Figure 1. General outline of BBQ core /transport iteration

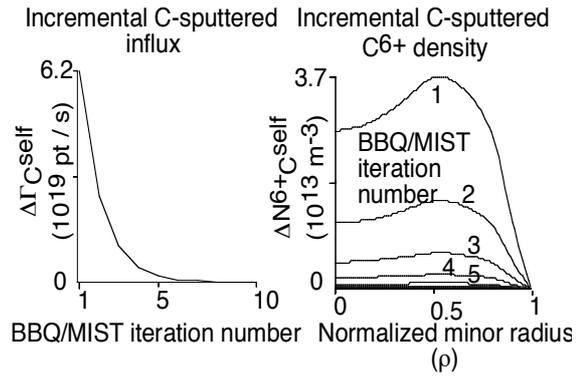


Figure 2. Example of BBQ/MIST iteration. (left) incremental C influx, (right) incremental $N_{C^{6+}}$

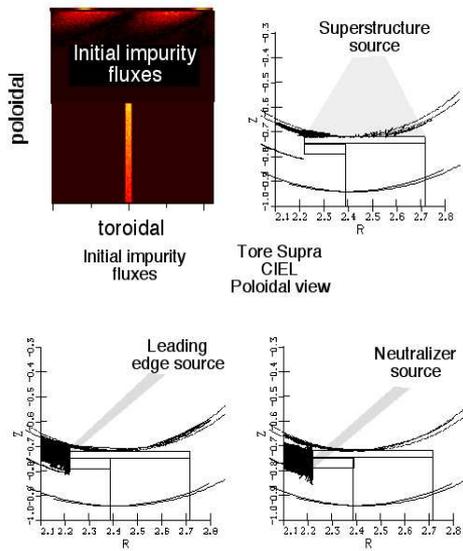


Figure 3. BBQ particles orbits: emitted flux distributions (top left) and particle orbits for superstructure, leading edge and neutralizer sources

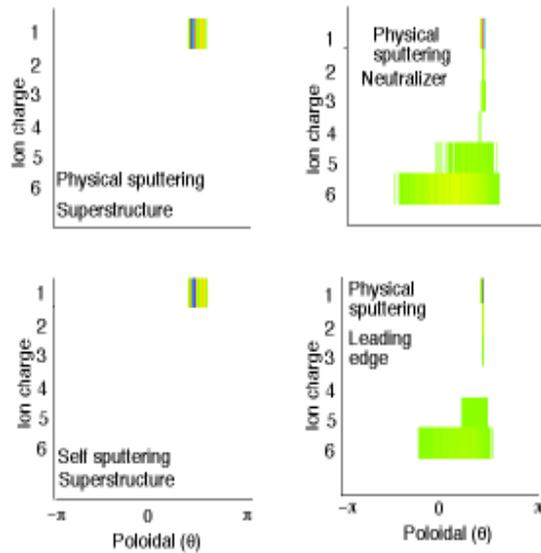


Figure 4. Charges state distributions of C entering core plasma influx (= color scale (rel.))

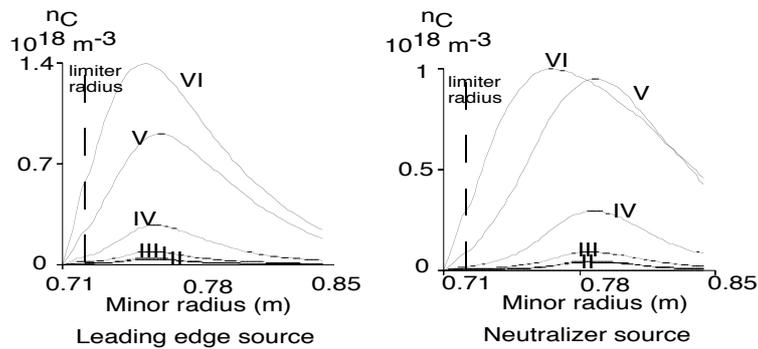


Figure 5. Radial distribution of C ion density for leading edge and neutralizer sources.