



# Properties of polycrystalline $\text{Hg}_{1-x}\text{Bi}_x\text{Ba}_2\text{Ca}_2\text{Cu}_3\text{O}_y$ superconductors

K.J. Song<sup>a,b</sup>, H.R. Khan<sup>c</sup>, H.J. Kim<sup>a,b</sup>, J.R. Thompson<sup>a,b,\*</sup>

<sup>a</sup>Oak Ridge National Laboratory, Oak Ridge, TN, USA

<sup>b</sup>Department of Physics, University of Tennessee, Knoxville, TN, USA

<sup>c</sup>FEM-Materials Physics Department, Katharinenstr. 17, D-73525 Schwäbisch Gmünd, Germany

## Abstract

Polycrystalline  $\text{Hg}_{1-x}\text{Bi}_x\text{Ba}_2\text{Ca}_2\text{Cu}_3\text{O}_y$  materials, with Bi-content  $x = 0.1, 0.2$  and  $0.3$ , were prepared by a direct, one-step synthesis. Characterization methods included X-ray diffraction, SEM, and studies of the magnetization  $M$ . The  $T_c$ 's were 133.5, 133 and 128 K, respectively. The persistent current density  $J_p$  falls off roughly exponentially with field  $H$  and temperature  $T$ . The irreversibility fields  $H_{\text{irr}}(T) \sim (1 - T/T_c)^n$ , with  $n \approx 3$ , were similar for the three materials. Analysis of the equilibrium  $M$  with London theory provided the London penetration depth  $\lambda(T)$ . © 2000 Elsevier Science B.V. All rights reserved.

**Keywords:** Magnetization; Penetration depth

The Hg-cuprate  $\text{HgBa}_2\text{Ca}_2\text{Cu}_3\text{O}_y$ , with three adjacent Cu–O layers has the highest known superconducting transition temperature  $T_c = 135$  K. With this high- $T_c$  and relatively simple tetragonal structure, the material is a natural candidate for further development. To improve the formation and stability of Hg-1223, partial substitutions of Hg by Bi, Pb, and Bi–Pb have been studied [1].

Here we investigate the properties of polycrystalline  $\text{Hg}_{1-x}\text{Bi}_x\text{Ba}_2\text{Ca}_2\text{Cu}_3\text{O}_y$  materials with Bi-content  $x = 0.1, 0.2$  and  $0.3$ , which were prepared directly from the metal oxides without use of a precursor [2]. Using a SQUID-based MPMS-7 magnetometer, the isothermal magnetization  $M(H)$  was measured at temperatures  $T$  between 5 and 120 K in magnetic fields  $H$  up to 6.5 T. Values of  $M$  were corrected for the paramagnetic background, measured for  $T$  up to 295 K and extrapolated to low temperatures.

Measurement of the low-field susceptibility  $4\pi\chi$  in  $H = 4$  Oe revealed  $T_c$ 's of 133.5, 133 and 128 K, respectively. At 5 K, the zero-field-cooled (ZFC) values of  $4\pi\chi$

were  $-0.80$ ,  $-0.67$  and  $-0.13$ ; the corresponding field-cooled (FC or Meissner) signals were  $4\pi\chi = -0.50$ ,  $-0.28$  and  $-0.04$ , respectively. These values are lower bounds for the superconducting volume fraction. For  $x = 0.1$  and  $0.2$ , the ZFC and FC transitions have “two-step” structures due to decoupling of current flow between grains. With Bi-content  $x = 0.3$ , there was only a single “one-step” transition.

The persistent current density  $J_p$  was obtained from the irreversible magnetization using the Bean critical state model. At low temperatures, the  $J_p$  values for the  $x = 0.1$  and  $0.2$  materials almost coincide. They begin to separate near 40 and by 77 K, the  $J_p$  for the  $x = 0.2$  material is substantially larger than that for the  $x = 0.1$  sample. With  $x = 0.3$ , the current density is considerably smaller than either of the other materials, qualitatively consistent with a reduced superconducting volume fraction. At low and intermediate temperatures,  $J_p$  falls off quasi-exponentially with  $H$  and  $T$ , then drops more quickly as one approaches the irreversibility line. The irreversibility field  $H_{\text{irr}}(T)$  was obtained from the  $J_p(T)$  results using the criterion that  $J_p(T) = 10^3$  A/cm<sup>2</sup>; at this level, the isothermal  $J$  has decreased by a factor of  $10^3$  or more from its low- $T$  value and is falling rapidly. The results follow an approximate power-law dependence

\* Corresponding author. Oak Ridge National Laboratory, P.O. Box 2008, Oak Ridge, TN 37831-6061, USA.

E-mail address: jrt@utk.edu (J.R. Thompson)

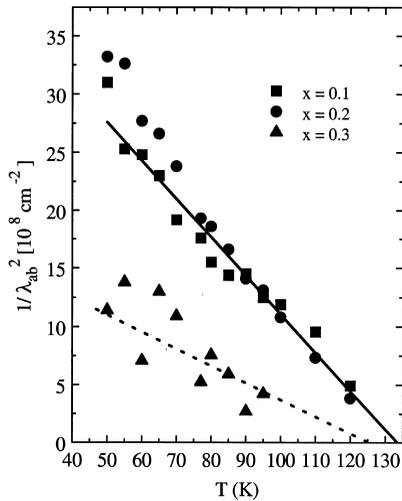


Fig. 1. The in-plane London penetration depth  $\lambda_{ab}$ , plotted as its inverse square, versus temperature  $T$ . Results are shown for  $\text{Hg}_{1-x}\text{Bi}_x\text{Ba}_2\text{Ca}_2\text{Cu}_3\text{O}_y$  with Bi-content  $x = 0.1, 0.2$  and  $0.3$ . Solid line is fitted to  $x = 0.1$  data and shows a Ginzburg–Landau dependence; dotted line is schematic.

$H_{\text{irr}}(T) \propto (1 - T/T_c)^n$  with similar exponents  $n \approx 3$  in each case.

Further studies of the equilibrium magnetization  $M$  were analyzed using standard London limit theory [3], which provides a logarithmic field dependence  $M \sim \ln(H)$ . The averaged magnetization  $\langle M \rangle$  at field  $H$  accurately followed this dependence. For a polycrystal, the slope  $dM/d \ln(H) = \phi_0/64\pi^2\lambda_{ab}^2$  [3]. Results for the London penetration depth in the  $ab$  plane,  $\lambda_{ab}(T)$ , are shown in Fig. 1 as a function of temperature. For the materials with  $x = 0.1$  and  $0.2$ , the resulting values are very similar. In the plot of  $1/\lambda_{ab}^2(T)$  versus  $T$ ,

Ginzburg–Landau theory provides a linear variation near  $T_c$ . This dependence is indeed observed, over an extended range of  $T$ ; note the solid line, fitted to all the data for  $x = 0.1$ , which intersects the axis very close to the observed low field  $T_c$ . A similar extended region of linear dependence was observed in a series of Hg-1201 materials [4]. Extrapolating linearly to  $T = 0$  gives  $\lambda_{ab}(0) = 150$  nm, similar to the values for Bi-free Hg-1223 prepared less directly [5]. For the material with  $x = 0.3$ , the data are more noisy and the dashed line schematically shows the Ginzburg–Landau dependence; no corrections have been made for a reduced volume fraction, which likely accounts for the larger *apparent* values of  $\lambda$  in this case.

### Acknowledgements

ORNL is managed by Lockheed Martin Energy Research Corp. for the USDOE under contract DE-AC05-96OR-22464. Work at the FEM was supported by Deutsche Forschungsgemeinschaft contract DFG-Khan 20/1-1.

### References

- [1] K. Isawa, A. Tokiwa-Yamamoto, M. Itoh, S. Adachi, H. Yamauchi, *Appl. Phys. Lett.* 65 (1994) 2105.
- [2] D. Loeblich, H.R. Khan, *Czech. J. Phys.* 46 (Suppl. S3) (1996) 1183.
- [3] V.G. Kogan, M.M. Fang, S. Mitra, *Phys. Rev. B* 38 (1988) 11958.
- [4] J.R. Thompson, H.R. Khan, K.J. Song, *Physica C* 272 (1996) 171.
- [5] J.R. Thompson, in: A.V. Narlikan (Ed.), *Studies of High Temperature Superconductors*, Vol. 26, Nova Science Publishers, Commack, NY, 1998, pp. 113–131.