

Overview of plasma-wall interactions in the SSPX spheromak experiment



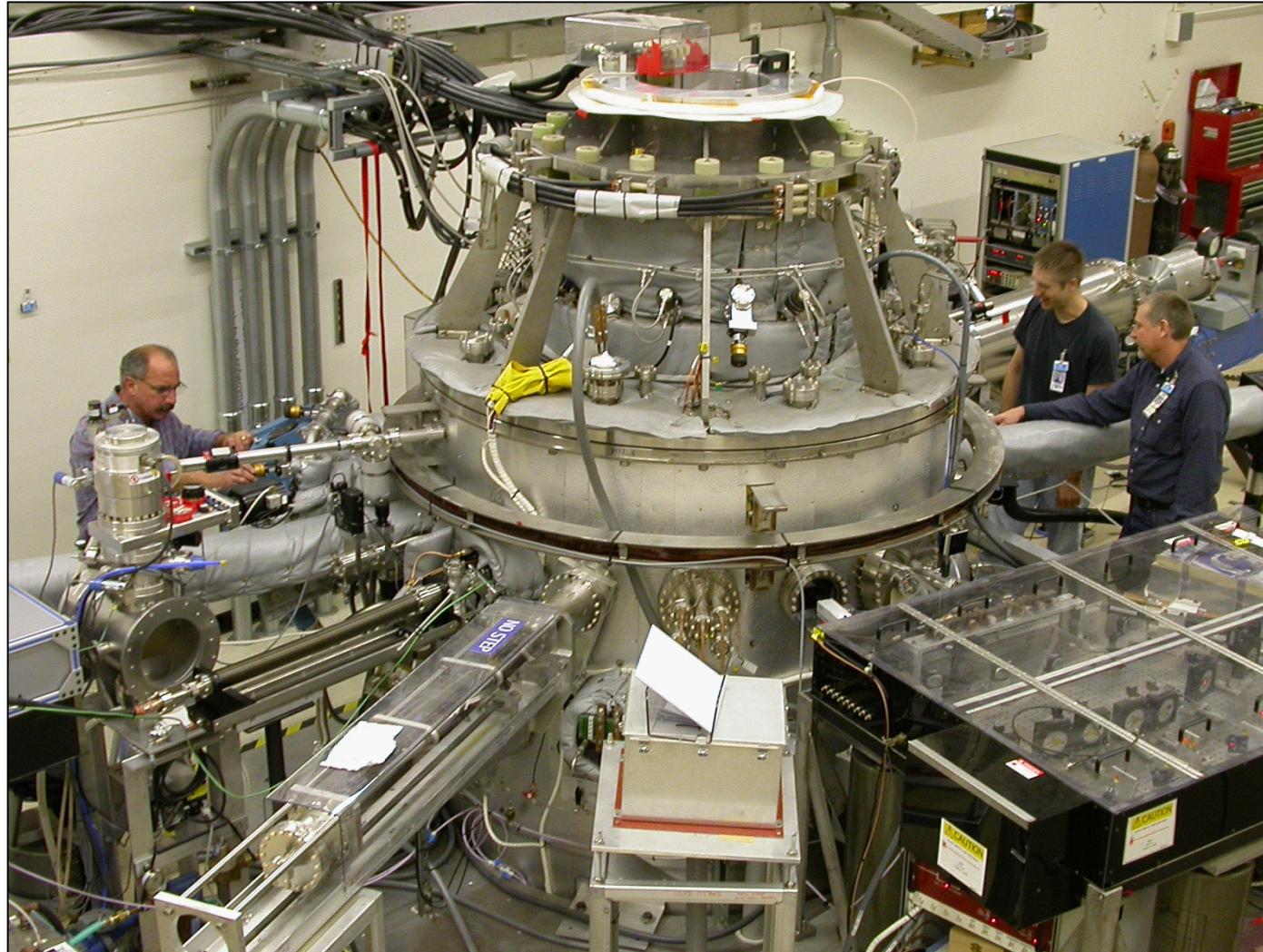
The SSPX Team: D.N. Hill, C. Holcomb, E.B. Hooper, H. McLean, C. Romero-Talamas, B.W. Stallard, R.D. Wood, and S. Woodruff

Lawrence Livermore National Laboratory

Outline of this talk

- Overview of the spheromak concept and purpose of the **Sustained Spheromak Physics eXperiment**
- Impurity control
- Power Balance
- Density control
- Future directions

The SSPX Spheromak began operating at LLNL in 1999

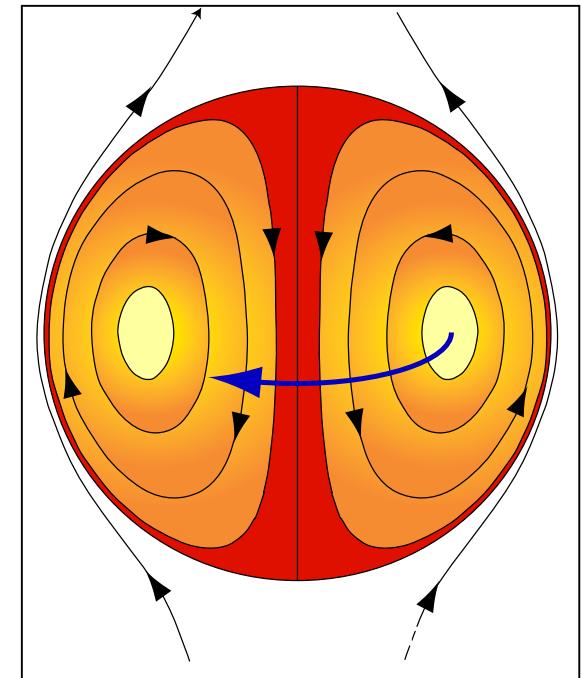


The SSPX spheromak



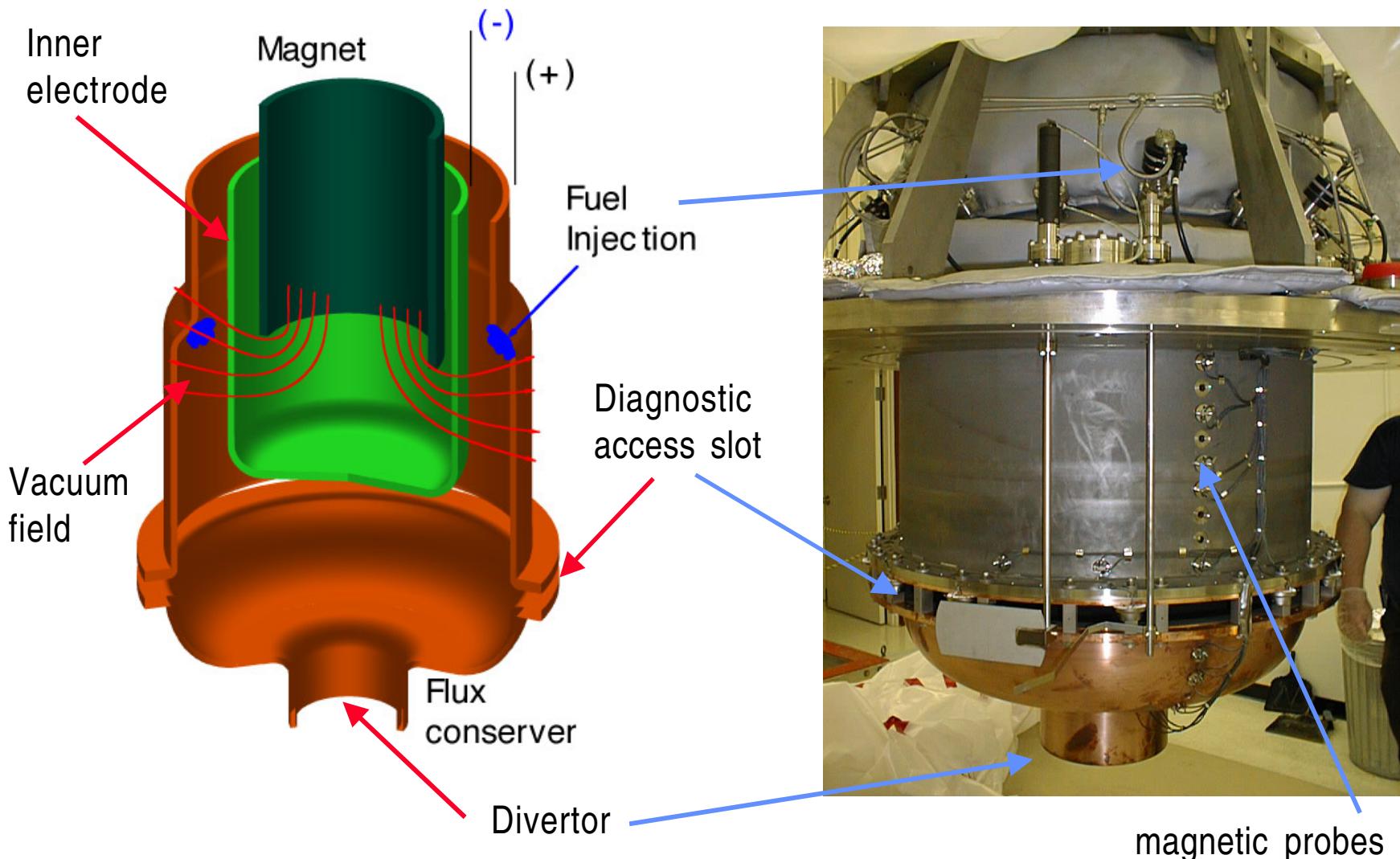
The Spheromak is a “self-organized” plasma configuration

- A low-aspect ratio ($R/a \gg 1$) toroidal plasma with
 - toroidal magnetic field on axis
 - poloidal magnetic field on edge
 - force-free currents parallel to fields
 $(\nabla P = j \times B \approx 0) \Rightarrow \nabla \times B = \lambda B \quad \lambda = j/B$
- Internal confining magnetic fields result from plasma currents in a self-consistent manner:
more efficient to generate fields in hot plasmas
- Magnetic fluctuations maintain configuration near that with minimum internal field energy:
a Taylor “relaxed state”.



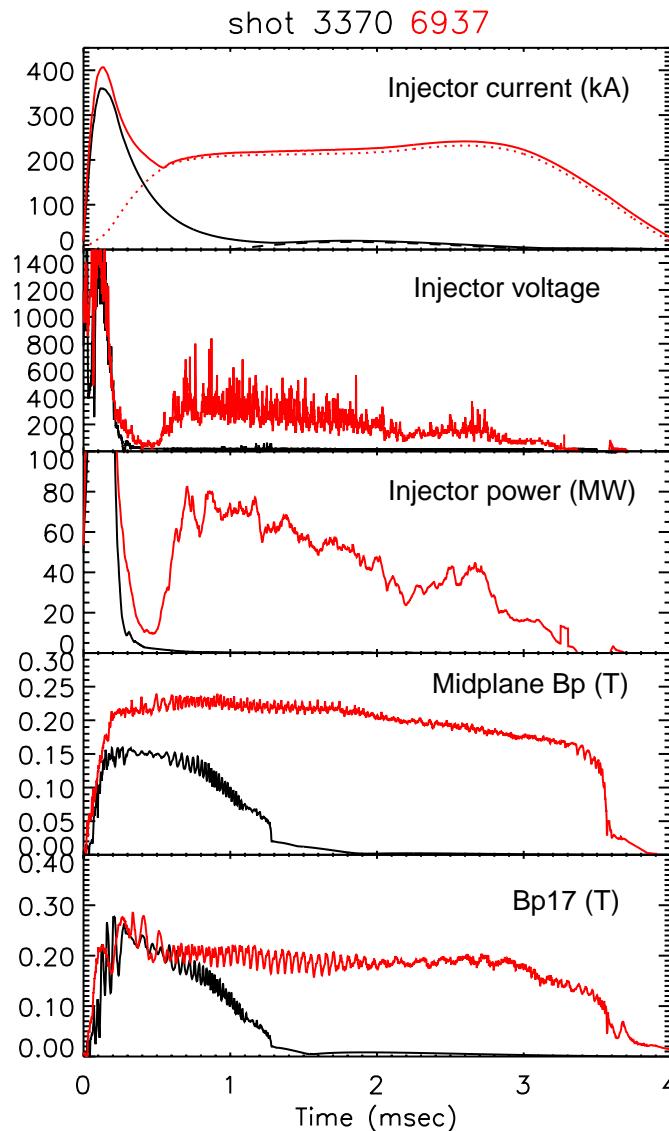
Self-generated magnetic fields point to potential for an attractive fusion reactor concept.

We use DC coaxial injection to form spheromak plasmas in SSPX

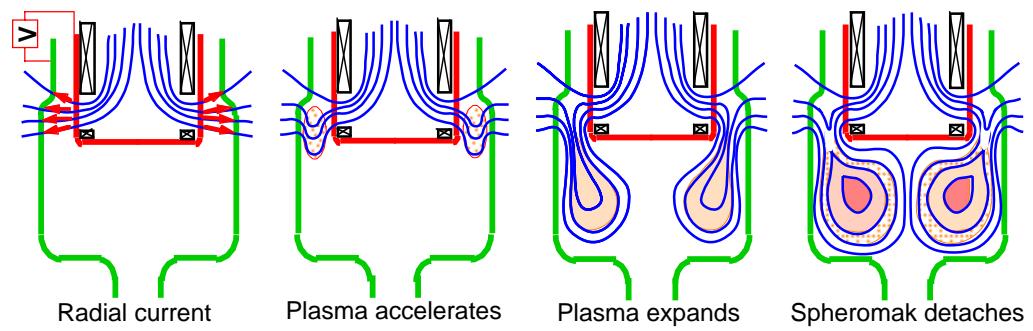




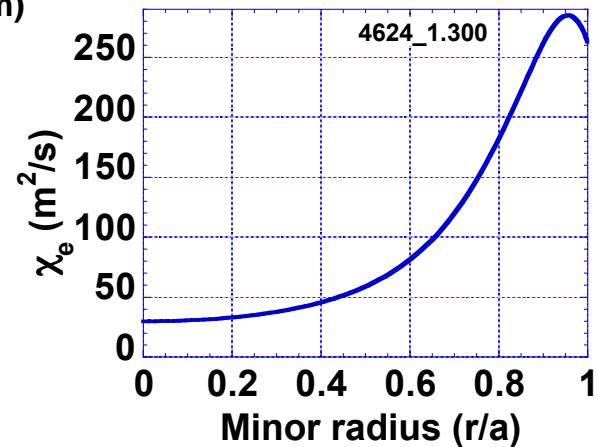
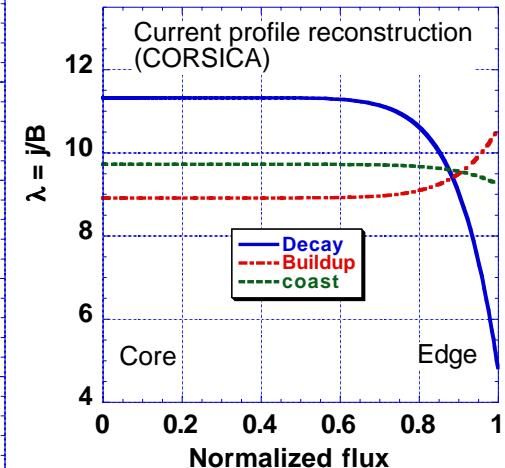
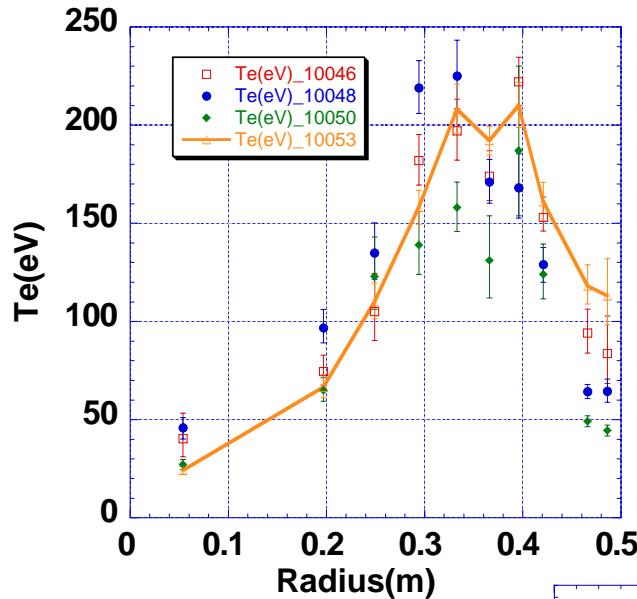
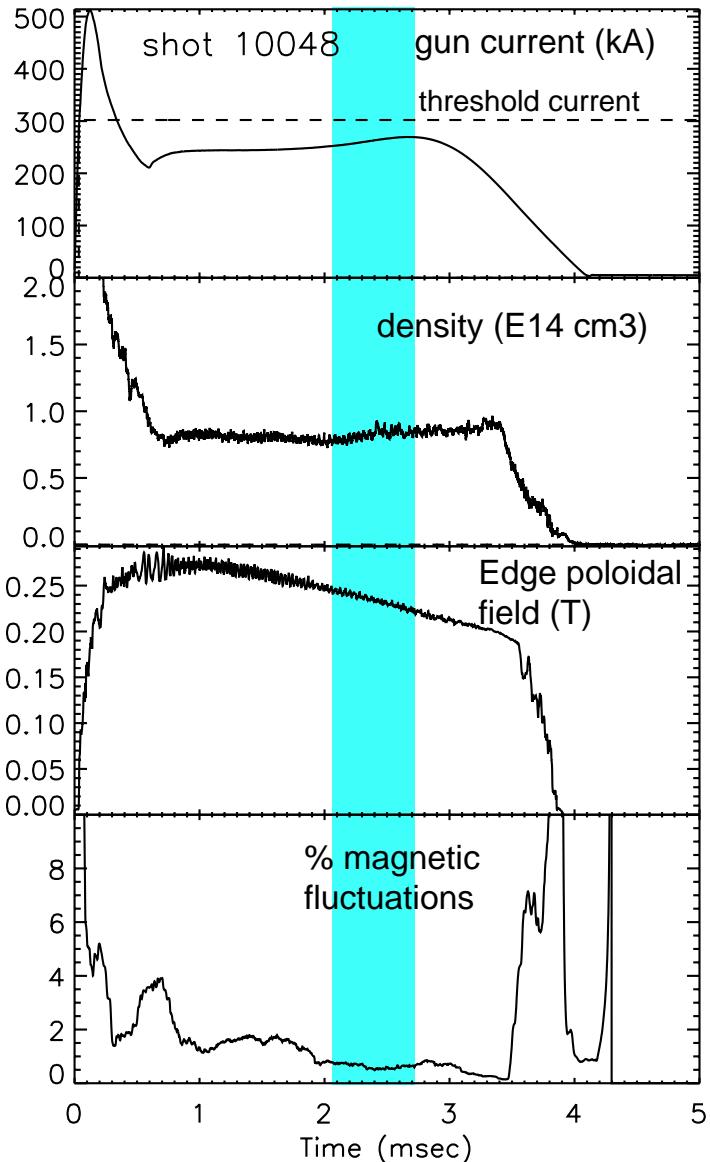
Standard formation and sustainment operation



- High current pulse forms spheromak (peak current well above threshold).
- Plasma is sustained with lower current from a second bank – current must remain above threshold.
- Peak edge poloidal magnetic field (and hence toroidal plasma current) is proportional to peak current.
- Threshold current depends on vacuum magnetic field geometry.



Fluctuation levels have been reduced by controlling the injection current

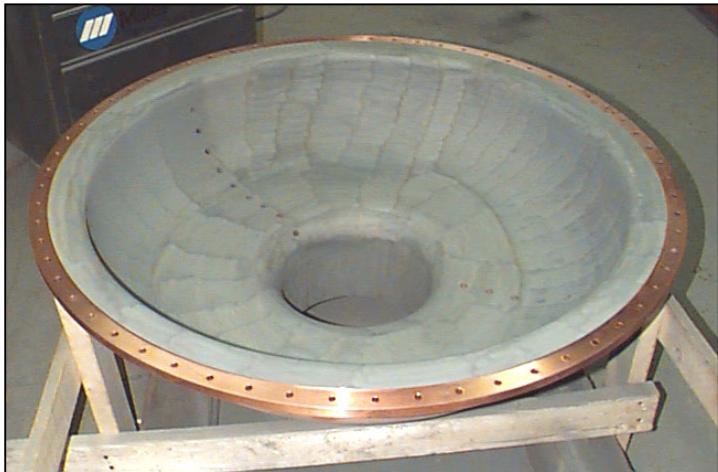


Lower drive current reduces edge current density and suppresses low-order modes.

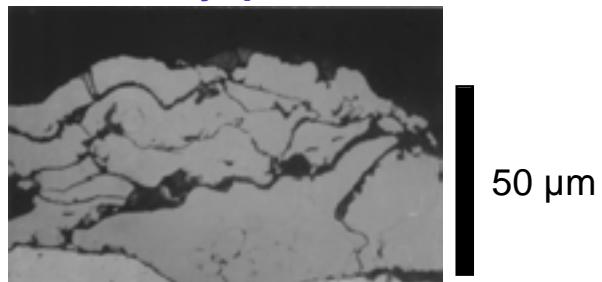
Proper wall-conditioning is required to obtain plasmas with low Z_{eff}



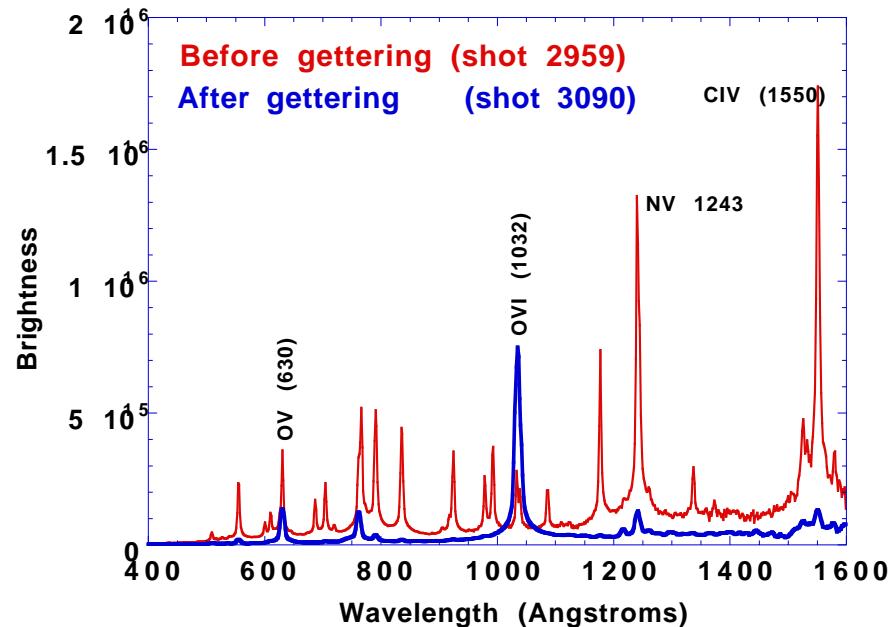
Plasma-sprayed W coatings



Relatively porous surface



D. Buchenauer, SNL

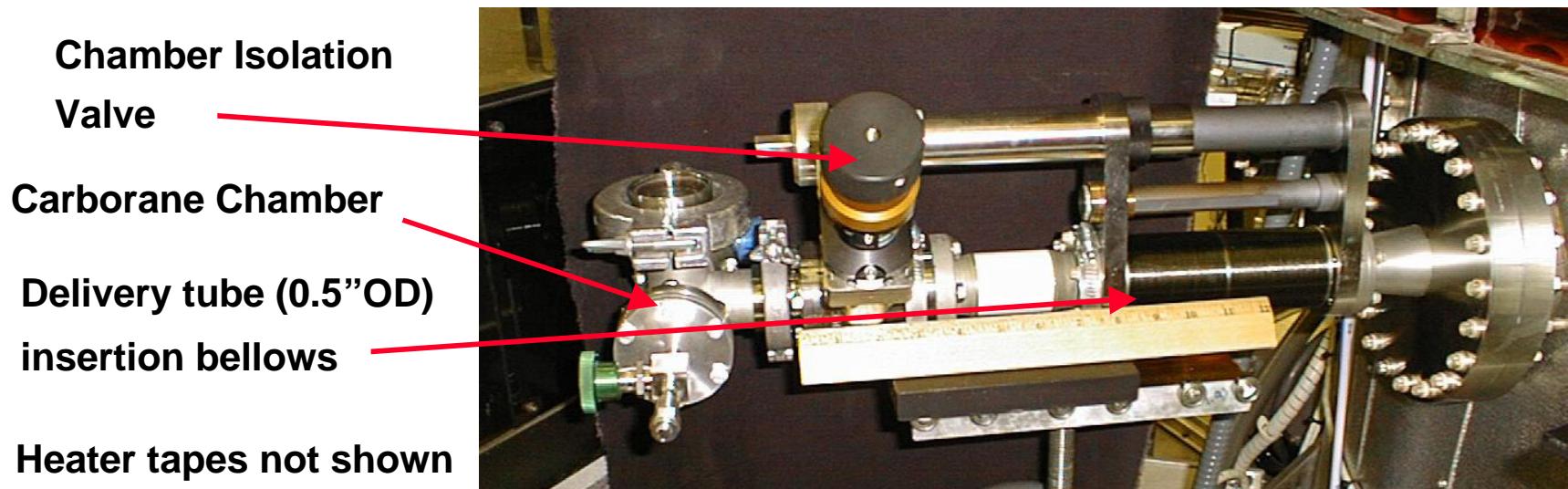


- Impurity radiation lowers T_e and maximum B field.
- Conditioning Processes:
 - high temperature bake (165 C)
 - hydrogen glow discharge cleaning
 - titanium gettering (every 4 shots)
 - helium plasma operation (12 shots)

Boronization using Carborane ($C_2B_{10}H_{12}$): Nontoxic, Non-explosive, and Inexpensive



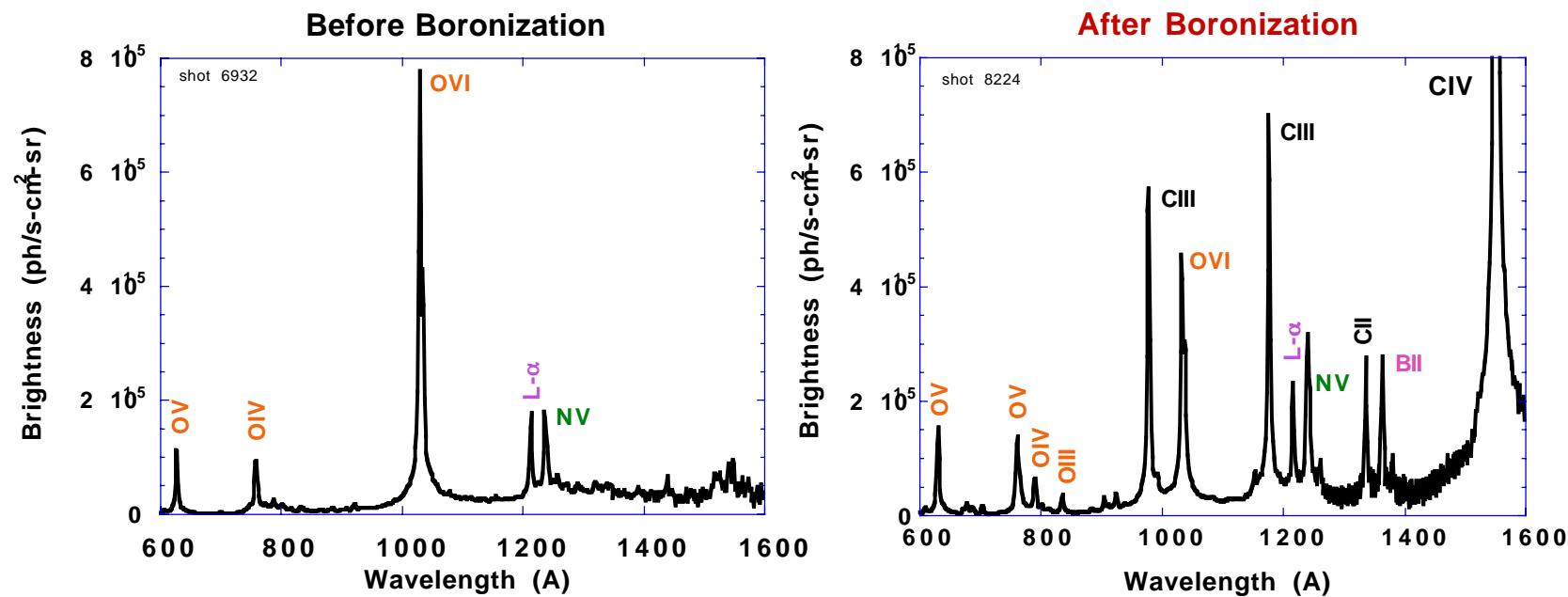
- Does not require special handling equipment
 - No special vacuum equipment, filters, etc.
 - Non-gaseous
- Nontoxic, non-explosive and simple hardware = inexpensive
- Boronization procedure:
 - He GDC and vessel baked to ~150C
 - GDC total pressure ~75 mTorr (90% helium 10% Carborane)
 - Vapor pressure rises x10 with each 10 C rise in carborane temp
 - Deposition rate ~40 nm/hr
 - Current density ~ 6 μ amps/cm²



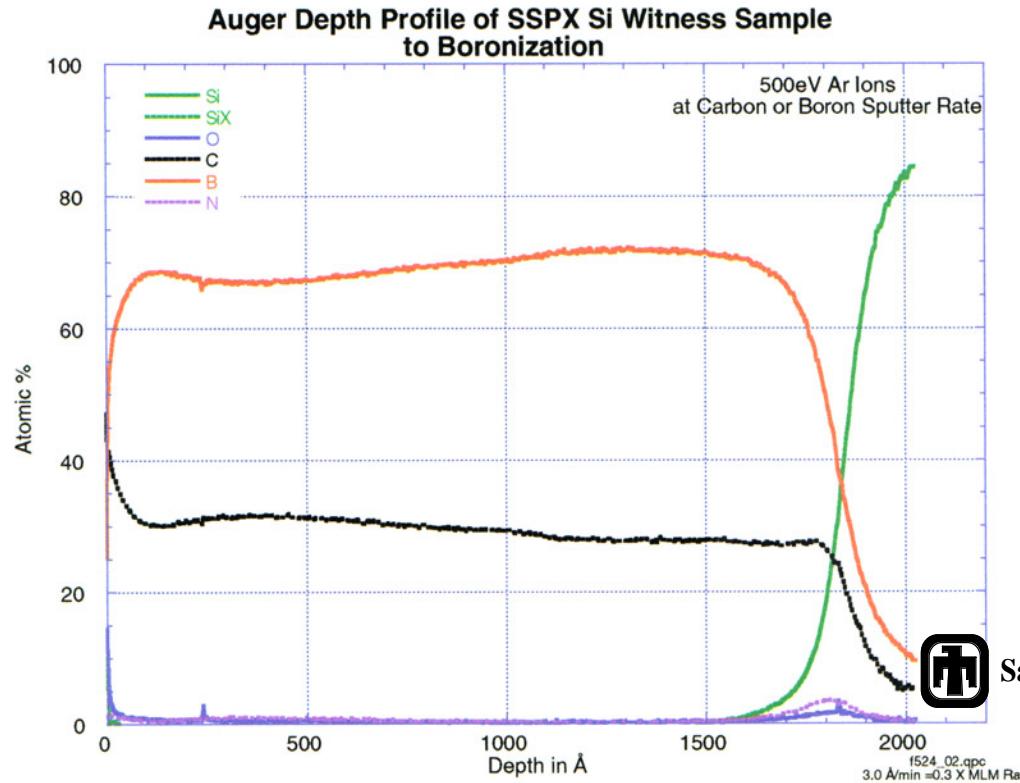
Increased Impurity Radiation Observed after Boronization with Carborane



- Deposition layer ~180 nm thick contains boron and carbon in a 70:30 ratio
- Increased carbon content with elevated electron density
- Oxygen concentration essentially unchanged
- After ~150 discharges:
 - Boron line radiation dropped to pre-boronization level
 - Carbon radiation dropped ~x40; yet x20 higher than before boronization



170nm boron–carbon film produced during 4hr helium GDC during 165 C bake



Depth calibration
based on boron
sputtering

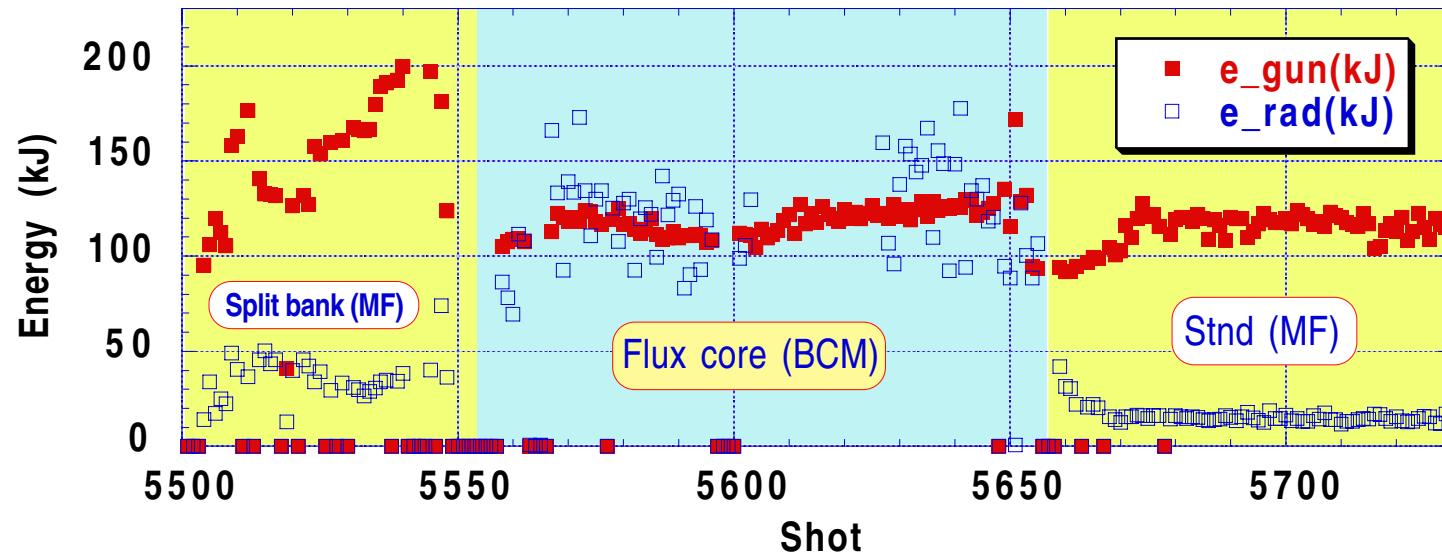
Dean Buchenauer

Sandia National Laboratories



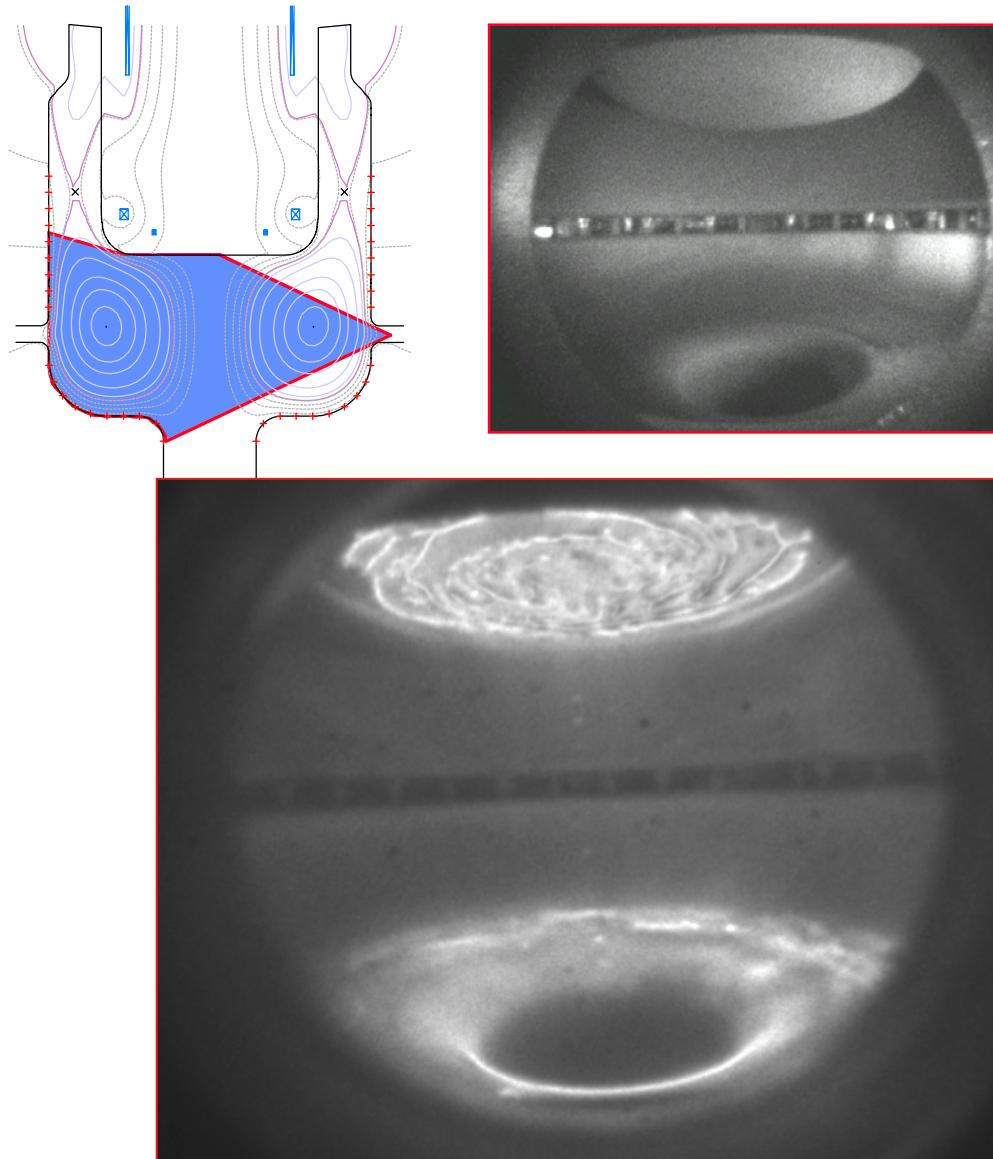
- Coupons located in diagnostic slot at flux conserver radius.
- Thin (50Å) layer with B:C=1.9:1 may be reflect shutdown conditions.
- Boron lines observed for about 100 discharges, carbon remains longer.

In clean plasmas, charge exchange and impurity radiation are a small fraction of input power



- Bolometer close to the plasma at the midplane provides integrated radiation and CX energy losses.
- Impurity radiation dominated by OVI (carbon is fully stripped).
- $P_{rad} \approx 15\%$ of P_{input} for clean plasmas.
- Melting of divertor plate in Flux core configuration produced almost 100% radiation loss!

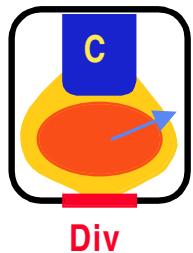
High speed imaging shows complex plasma-surface interactions



Images courtesy of
C. Romero-Talama (Caltech)

- **4 μ sec exposure time in visible light midway through discharge**
- **Bright spots on bottom of inner electrode points to possible filamentation on open field lines.**
- **Interaction with lower flux conserver shows both limiter and electrode characteristics.**

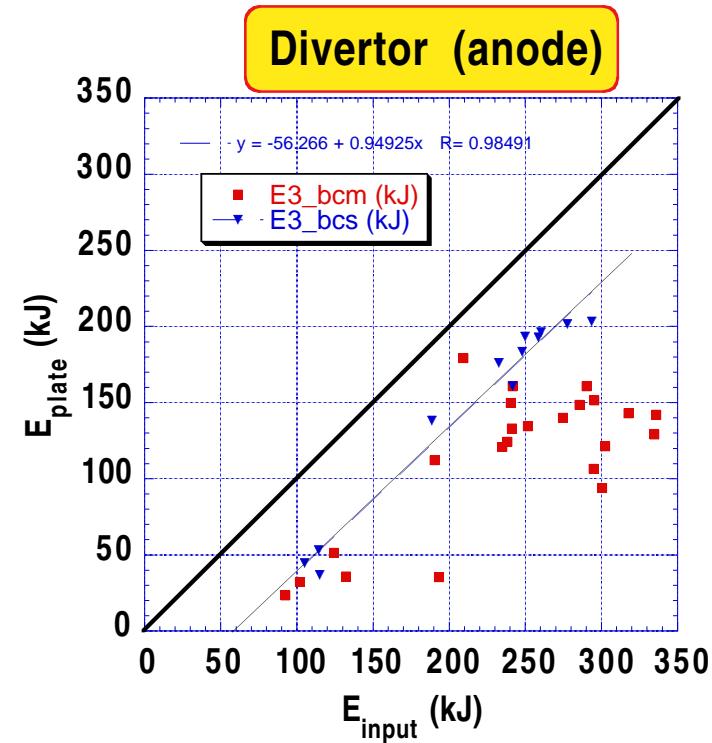
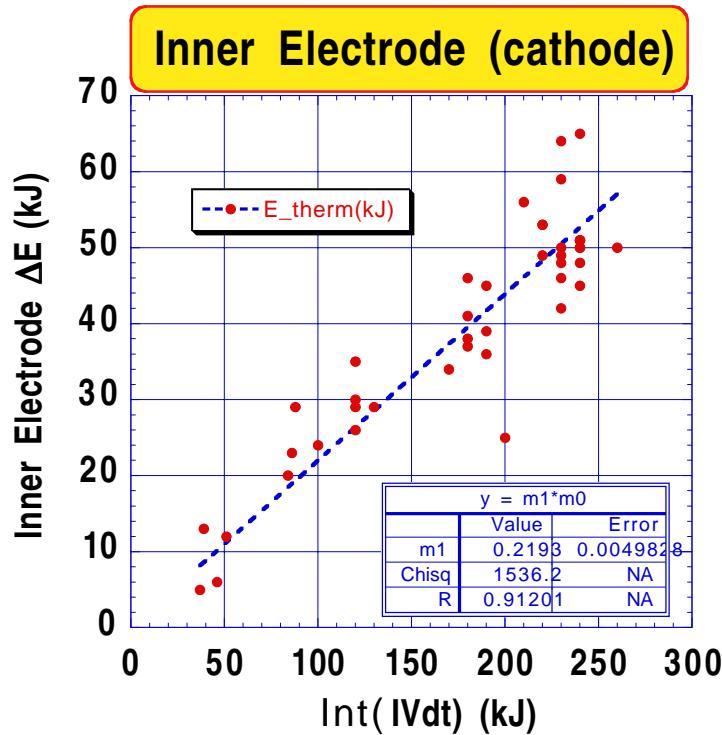
Power balance measurements show that most of the power goes to the divertor (discharge anode)



$$\begin{aligned}
 P_{\text{inj}} &= I_g V_g = P_{\text{cx}} + P_{\text{rad}} + P_{\text{cath}} + P_{\text{anode}} \\
 &= P_{\text{bolom}} + P_{\text{inr-elec}} + P_{\text{div}} \\
 1.0 &= 0.15 + 0.2 + 0.7
 \end{aligned}$$

$\cdot W_B + \cdot W_{nkT}$
 $\cdot W_B + \cdot W_{nkT}$
+ 0 (stationary)

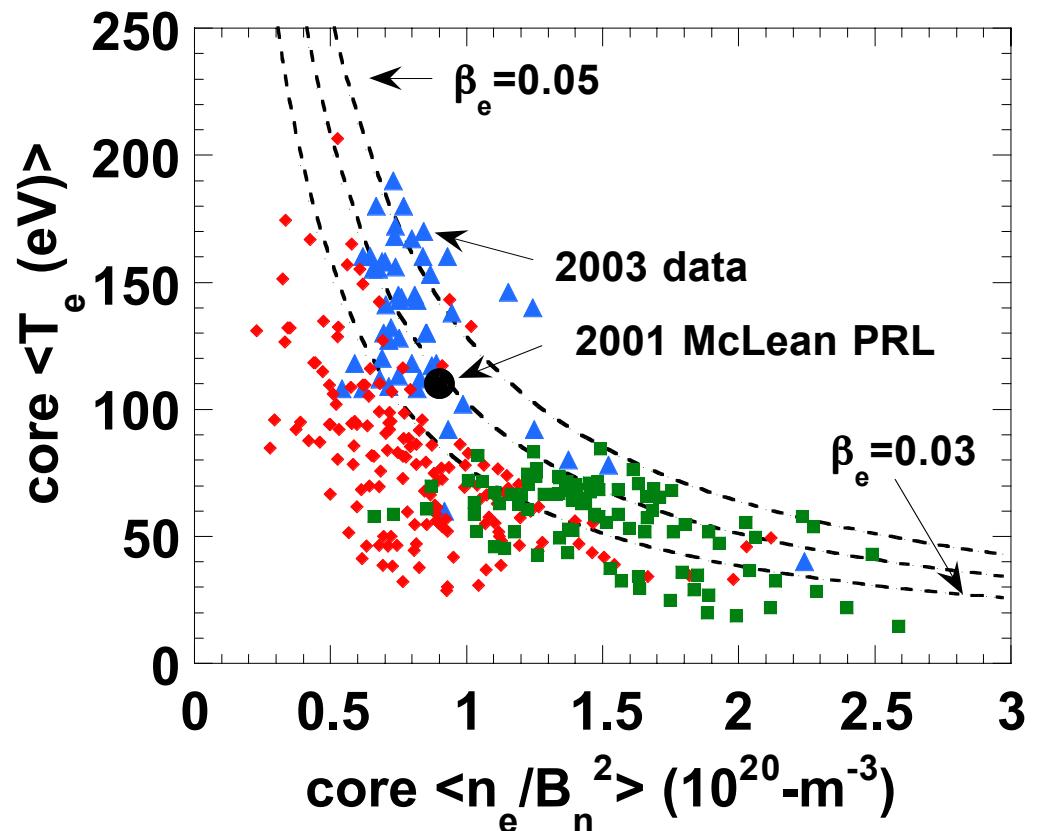
High level power balance



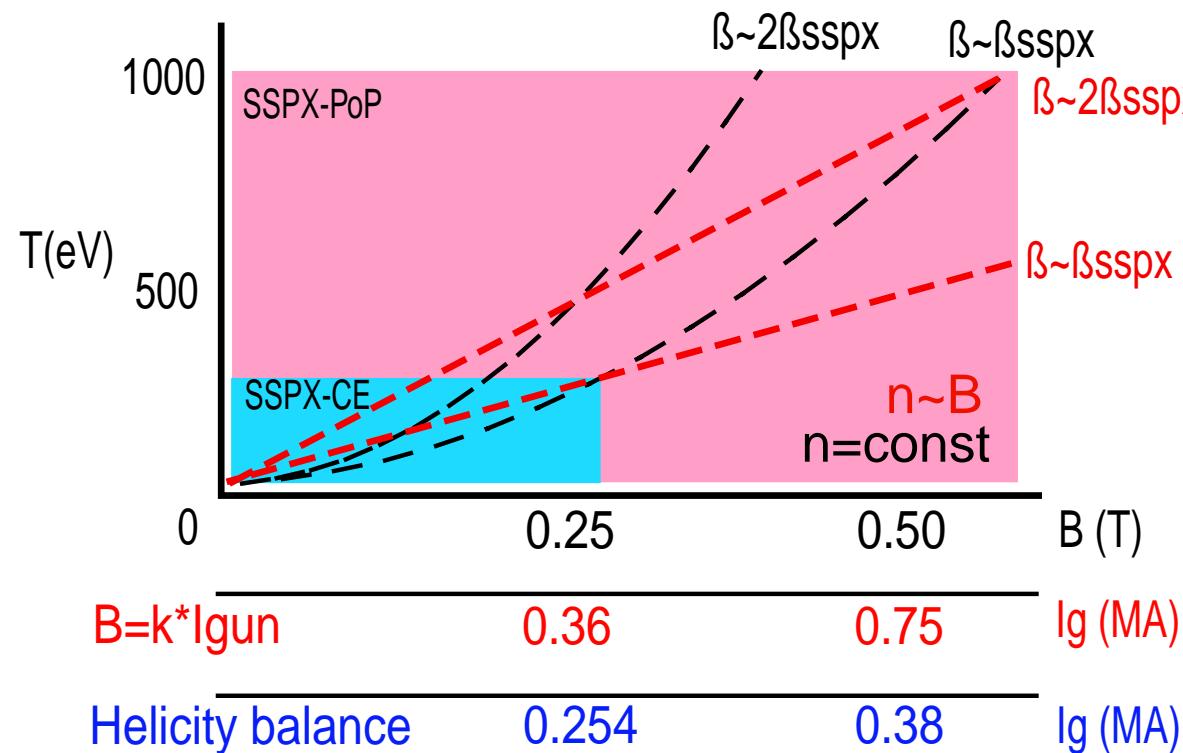
Temperature measurements point to importance of density control



- Core T_e measured in sustained plasmas.
- Upper limit observed for $\beta = 2\mu_0 n k T / B^2$.
- Higher fields or lower density should $\uparrow T_e$.
- In the spheromak, higher fields are obtained by higher plasma current.
- So far, I_{tor} and $B_{sphere} \propto I_{gun}$.

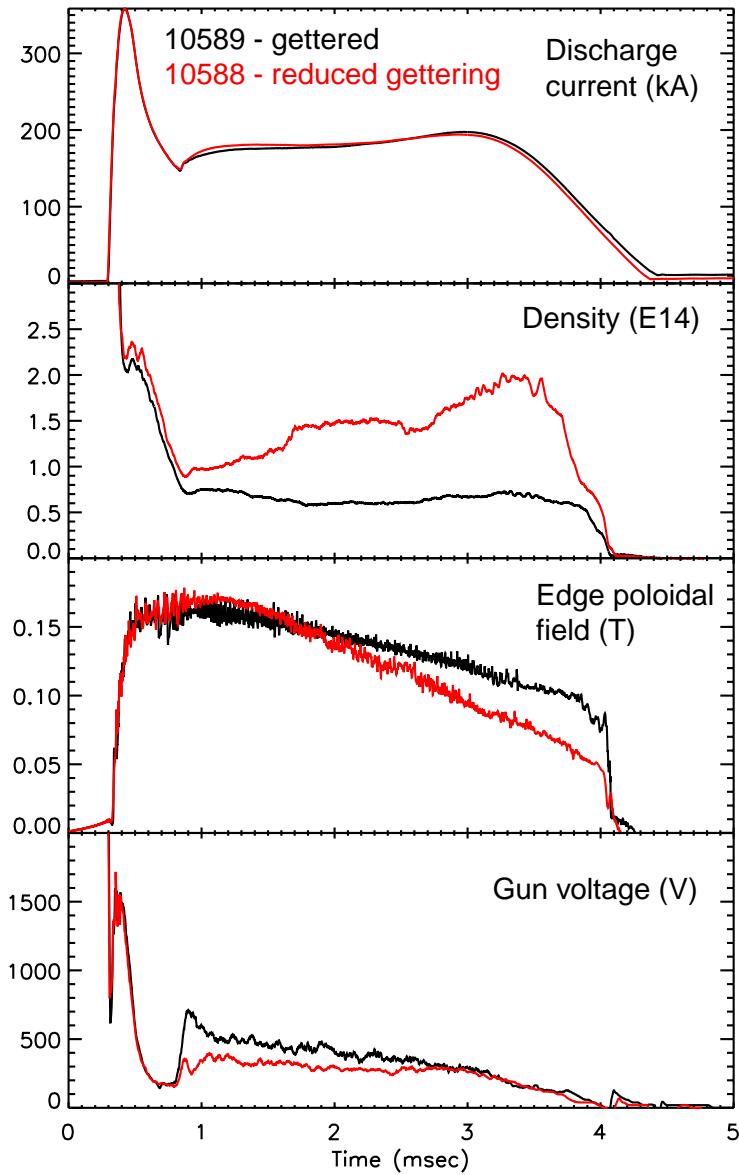


Extrapolation to next-step spheromaks point to the importance of density control



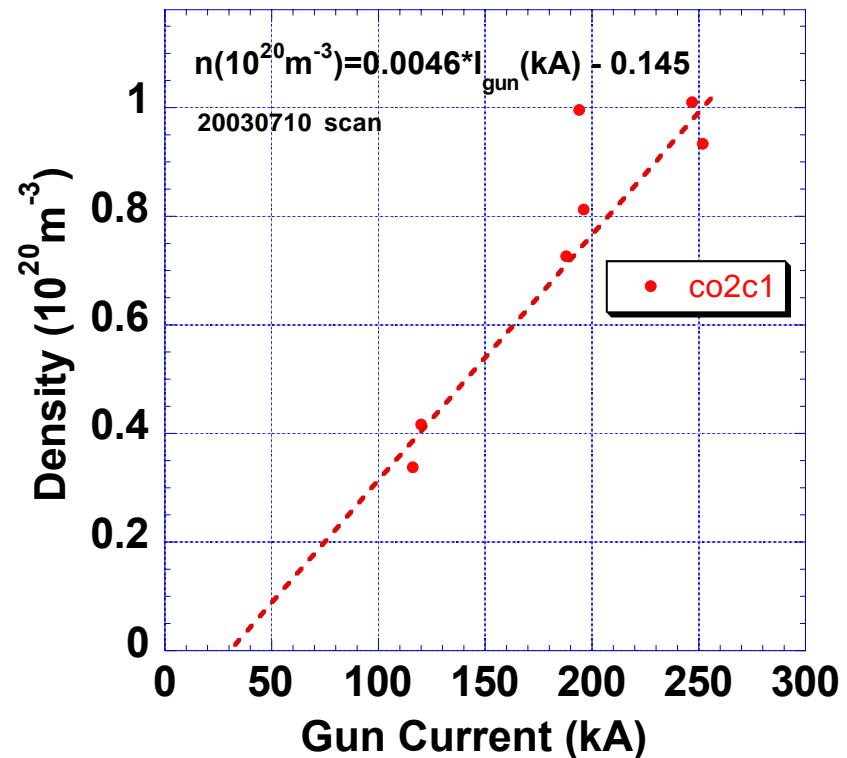
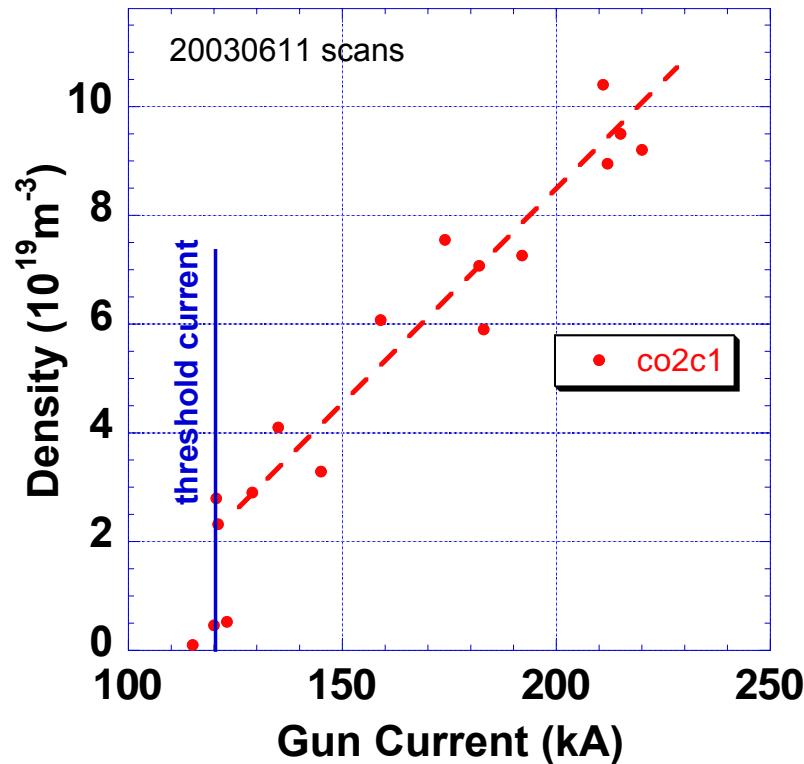
- Target 1 keV spheromak plasma. What are the device requirements?
- Transport scaling sets device size (minor radius).
- Beta limit and density scaling determine field and toroidal current.
- Field generation efficiency (B_p/I_{gun}) sets bank and injector requirements.

Spheromak density is sustained by recycling



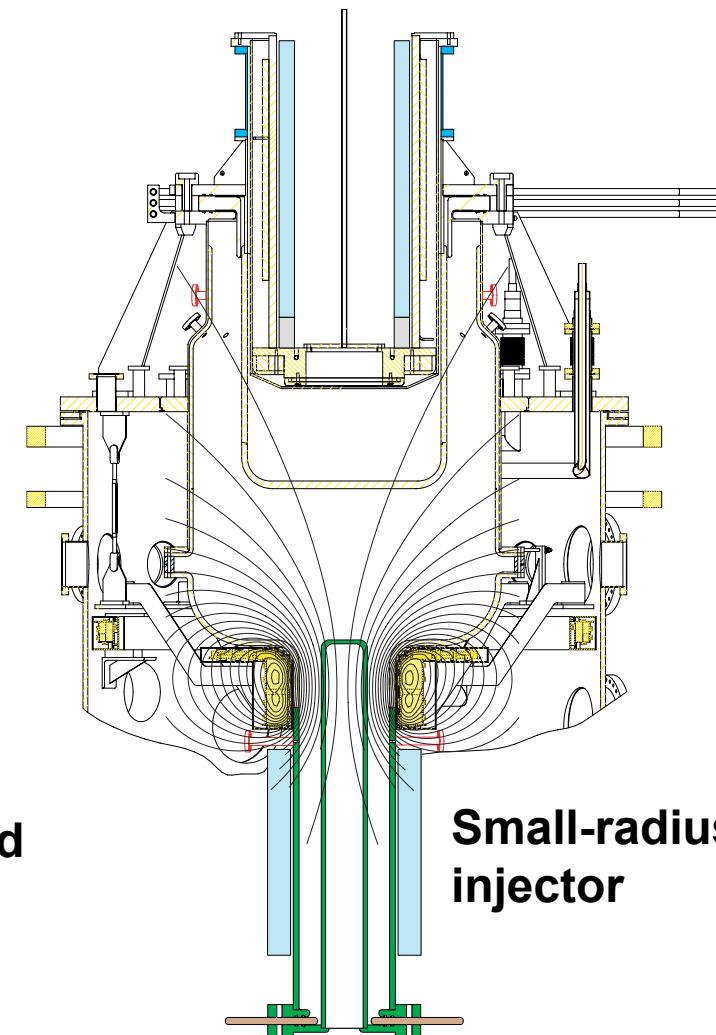
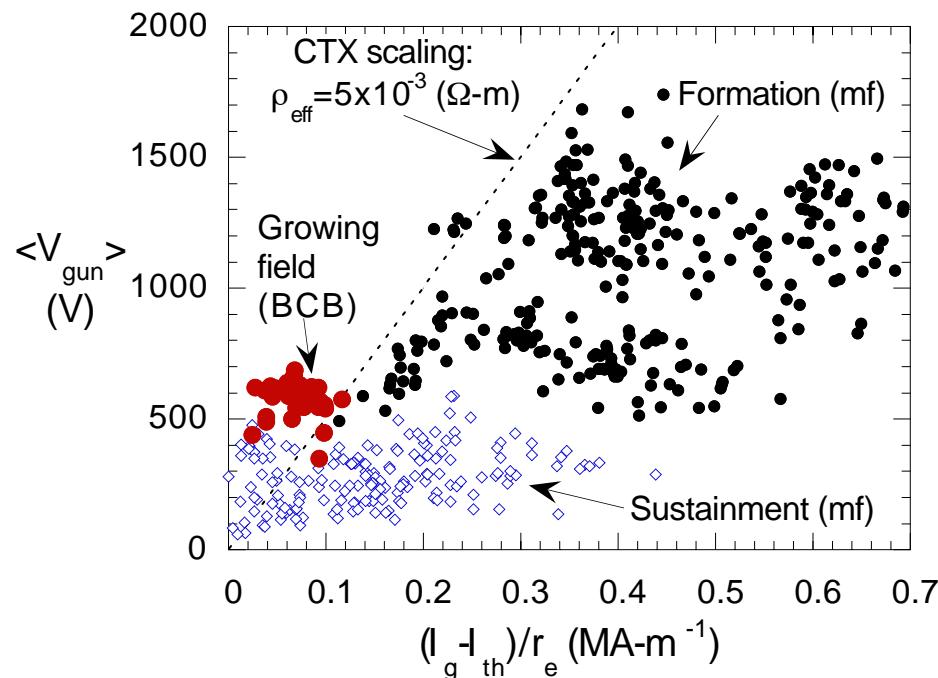
- Initial density corresponds to full ionization of gas puff @ -250 μ sec.
- Density decay time (~350 μ sec) represents ion loss to walls-much shorter than magnetic field decay time.
- Steady-state density depends on wall conditioning and gun current, not initial gas puff.
- Gettering in main chamber only, not in coaxial source region.
- Effectiveness of gettering reduced after 3 – 4 discharges.

Discharge current determines density



- Current scan at fixed flux explores effect of sustainment threshold.
- Current scan at fixed $\lambda = I_{\text{gun}} / \phi_{\text{gun}}$ explores scaling at optimal current.

Small radius gun should increase inductive voltage and increase the spheromak magnetic fields.



- Smaller gun increases discharge voltage and coaxial magnetic field
- Factor of two voltage increase expected
 - $V \sim L \sim \log(R_o/R_i)$
 - $L_{\text{new}} / L_{\text{old}} = .30 / .17$
 - Higher helicity input rate $\sim V_{\text{gun}} \Psi_{\text{gun}}$



Summary

- Spheromak potentially offers attractive reactor concept for magnetic fusion energy
- We are producing driven spheromaks with low-amplitude fluctuations and peaked temperature and pressure profiles
 - Power balance shows low radiation losses
 - $T_{e0} \geq 200\text{eV}$
 - Wall conditioning key to obtaining best performance
- Magnetic field generation and density control are important aspects to showing favorable scaling to next step devices
 - Attractive reactor concept depends on efficient field generation
 - Optimal density avoids beta limits and gives high T_e .
- A new coaxial injector will be installed to increase the magnetic field.